

A Spin 1/2 Particle with Three Additional Characteristics in Presence of External Fields, General Theory within the Gel'fand–Yaglom Formalism

V.V. Kisel,^{*} E.M. Ovsyuk,[†] P.O. Sachenok,[‡] and A.S. Martynenko[§]
Mozyr State Pedagogical University named after I.P. Shamyakin

A.V. Bury[§] and V.M. Red'kov[¶]
*B. I. Stepanov Institute of Physics of NAS of Belarus,
 68 Nezavisimosti Ave., 220072 Minsk, BELARUS*

Within the general method by Gel'fand – Yaglom, starting with the extended 28-component representation of the Lorentz group (it includes four bispinors and one spinor of the 3rd rank), for a spin 1/2 particle, we construct a relativistic P -invariant generalized system of the first order equations for a spin 1/2 particle with three additional to electric charge characteristics.

First, the model is developed for a free particle, and the system of spinor equations is derived; then it is transformed to spin-tensor form. In this spin-tensor form, we take into account the presence of external electromagnetic fields. After eliminating the accessory variables of the complete wave function, we derive the minimal 4-component Dirac-like equation, the last includes several new interaction terms which are interpreted as related to some additional electromagnetic characteristics of a spin 1/2 particle.

PACS numbers: 02.30.Gp, 02.40.Ky, 03.65.Ge, 04.62.+v

Keywords: Spin 1/2, Gel'fand – Yaglom, extended wave equations, additional characteristics, external electromagnetic field, Riemannian geometry.

1. Gel'fand – Yaglom formalism, and a new Dirac-like equation

We will construct a generalized relativistic equation for a particle with spin 1/2, using the extended set of irreducible representations of the proper Lorentz group

$$T = 4[(0, 1/2) \oplus (1/2, 0)] \oplus [(1/2, 1) \oplus (1, 1/2)]; \quad (1)$$

with the linking scheme

$$\begin{array}{ccc} 4(0, 1/2) & - & 4(1/2, 0) \\ | & & | \\ (1/2, 1) & - & (1, 1/2) \end{array} \quad (2)$$

First we will construct a matrix equation for a free particle ¹

$$(\Gamma_\mu \partial_\mu + M)\Psi = 0, \quad \mu = 1, 2, 3, 4. \quad (3)$$

In modified Gel'fand – Yaglom basis [2], the components of the complete wave function are determined as follows

$$\begin{aligned} \Psi^{\text{G-Y(mod)}} = & \left\{ (\Psi^{(1)})_{1/2,1/2}^{(0,1/2)}, (\Psi^{(1)})_{1/2,-1/2}^{(0,1/2)}, (\Psi^{(1)})_{1/2,1/2}^{(1/2,0)}, (\Psi^{(1)})_{1/2,-1/2}^{(1/2,0)}; \right. \\ & (\Psi^{(2)})_{1/2,1/2}^{(0,1/2)}, (\Psi^{(2)})_{1/2,-1/2}^{(0,1/2)}, (\Psi^{(2)})_{1/2,1/2}^{(1/2,0)}, (\Psi^{(2)})_{1/2,-1/2}^{(1/2,0)}; \\ & (\Psi^{(3)})_{1/2,1/2}^{(0,1/2)}, (\Psi^{(3)})_{1/2,-1/2}^{(0,1/2)}, (\Psi^{(3)})_{1/2,1/2}^{(1/2,0)}, (\Psi^{(3)})_{1/2,-1/2}^{(1/2,0)}; \\ & \left. (\Psi^{(4)})_{1/2,1/2}^{(0,1/2)}, (\Psi^{(4)})_{1/2,-1/2}^{(0,1/2)}, (\Psi^{(4)})_{1/2,1/2}^{(1/2,0)}, (\Psi^{(4)})_{1/2,-1/2}^{(1/2,0)} \right\} \end{aligned}$$

^{*}E-mail: vasiliy-bspu@mail.ru

[†]E-mail: e.ovsiyuk@mail.ru

[‡]E-mail: polinasachenok@gmail.com

[§]E-mail: anton.buryy.97@mail.ru

[¶]E-mail: v.redkov@ifanbel.bas-net.by

¹ Applying the *ict*-metric.

$$\begin{aligned} & \Psi_{1/2,1/2}^{(1,1/2)}, \Psi_{1/2,-1/2}^{(1,1/2)}, \Psi_{1/2,1/2}^{(1/2,1)}, \Psi_{1/2,-1/2}^{(1/2,1)}; \\ & \Psi_{3/2,3/2}^{(1,1/2)}, \Psi_{3/2,-3/2}^{(1,1/2)}, \Psi_{3/2,3/2}^{(1/2,1)}, \Psi_{3/2,-3/2}^{(1/2,1)}; \\ & \Psi_{3/2,1/2}^{(1,1/2)}, \Psi_{3/2,-1/2}^{(1,1/2)}, \Psi_{3/2,1/2}^{(1/2,1)}, \Psi_{3/2,-1/2}^{(1/2,1)} \}^T; \end{aligned} \quad (4)$$

the symbol T designates the matrix transposition; the matrix Γ_4 may be presented in the form

$$\Gamma_4 = \begin{vmatrix} c^{(1/2)} \otimes \gamma_4 & 0 \\ 0 & c^{(3/2)} \otimes I_2 \otimes \gamma_4 \end{vmatrix}, \quad (5)$$

where the spin blocks $c^{(1/2)}, c^{(3/2)}$ have the structure

$$c^{(1/2)} = \begin{vmatrix} c_{11'}^{(1/2)} & c_{12'}^{(1/2)} & c_{13'}^{(1/2)} & c_{14'}^{(1/2)} & c_{15'}^{(1/2)} \\ c_{21'}^{(1/2)} & c_{22'}^{(1/2)} & c_{23'}^{(1/2)} & c_{24'}^{(1/2)} & c_{25'}^{(1/2)} \\ c_{31'}^{(1/2)} & c_{32'}^{(1/2)} & c_{33'}^{(1/2)} & c_{34'}^{(1/2)} & c_{35'}^{(1/2)} \\ c_{41'}^{(1/2)} & c_{42'}^{(1/2)} & c_{43'}^{(1/2)} & c_{44'}^{(1/2)} & c_{45'}^{(1/2)} \\ c_{51'}^{(1/2)} & c_{52'}^{(1/2)} & c_{53'}^{(1/2)} & c_{54'}^{(1/2)} & c_{55'}^{(1/2)} \end{vmatrix}, \quad C^{(3/2)} = C_{55'}^{(3/2)}, \quad \gamma_4 = \begin{vmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{vmatrix}; \quad (6)$$

I_2 is a 2×2 unit matrix; the involved irreducible representations of the proper Lorentz group are numerated as follows

$$1, 2, 3, 4 \implies (0, 1/2), \quad 1', 2', 3', 4' \implies (1/2, 0), \quad 5 \implies (1, 1/2), \quad 5' \implies (1/2, 1).$$

According to [1, 2], we have

$$c^{(1/2)} = \begin{vmatrix} c_{11'} & c_{12'} & c_{13'} & c_{14'} & i\sqrt{3}c_{15'} \\ c_{21'} & c_{22'} & c_{23'} & c_{24'} & i\sqrt{3}c_{25'} \\ c_{31'} & c_{32'} & c_{33'} & c_{34'} & i\sqrt{3}c_{35'} \\ c_{41'} & c_{42'} & c_{43'} & c_{44'} & i\sqrt{3}c_{45'} \\ i\sqrt{3}c_{51'} & i\sqrt{3}c_{52'} & i\sqrt{3}c_{53'} & i\sqrt{3}c_{54'} & c_{55'} \end{vmatrix}, \quad C^{(3/2)} = 2C_{55'}, \quad (7)$$

where $c_{ij'}$ are some numerical parameters, on which special constraints will be imposed later on.

In order to simplify the calculations bellow, let us take into account the possibility to break the linking between repeated representations [1]; this is achieved by vanishing the corresponding parameters in the block $c^{(1/2)}$:

$$c^{(1/2)} = \begin{vmatrix} c_{11'} & 0 & 0 & 0 & i\sqrt{3}c_{15'} \\ 0 & c_{22'} & 0 & 0 & i\sqrt{3}c_{25'} \\ 0 & 0 & c_{33'} & 0 & i\sqrt{3}c_{35'} \\ 0 & 0 & 0 & c_{44'} & i\sqrt{3}c_{45'} \\ i\sqrt{3}c_{51'} & i\sqrt{3}c_{52'} & i\sqrt{3}c_{53'} & i\sqrt{3}c_{54'} & c_{55'} \end{vmatrix}, \quad c^{(3/2)} = 2c_{55'}. \quad (8)$$

Because we assume the absence states with spin 3/2, we should set $c_{55'} = 0$; then we arrive at

$$c^{(1/2)} = \begin{vmatrix} c_{11'} & 0 & 0 & 0 & i\sqrt{3}c_{15'} \\ 0 & c_{22'} & 0 & 0 & i\sqrt{3}c_{25'} \\ 0 & 0 & c_{33'} & 0 & i\sqrt{3}c_{35'} \\ 0 & 0 & 0 & c_{44'} & i\sqrt{3}c_{45'} \\ i\sqrt{3}c_{51'} & i\sqrt{3}c_{52'} & i\sqrt{3}c_{53'} & i\sqrt{3}c_{54'} & 0 \end{vmatrix}, \quad c^{(3/2)} = 0. \quad (9)$$

It is convenient to introduce the new notations:

$$\lambda_1 = c_{11'}, \quad \lambda_2 = c_{22'}, \quad \lambda_3 = c_{33'}, \quad \lambda_4 = c_{44'}, \quad (10)$$

$$\begin{aligned} \frac{\beta_1}{\sqrt{2}} &= -ic_{15'}, & \frac{\beta_2}{\sqrt{2}} &= -ic_{25'}, & \frac{\beta_3}{\sqrt{2}} &= -ic_{35'}, & \frac{\beta_4}{\sqrt{2}} &= -ic_{45'}, \\ \frac{\beta_5}{\sqrt{2}} &= ic_{51'}, & \frac{\beta_6}{\sqrt{2}} &= ic_{52'}, & \frac{\beta_7}{\sqrt{2}} &= ic_{53'}, & \frac{\beta_8}{\sqrt{2}} &= ic_{54'}; \end{aligned} \quad (11)$$

so we obtain

$$c^{(1/2)} = \begin{vmatrix} \lambda_1 & 0 & 0 & 0 & -\sqrt{\frac{3}{2}}\beta_1 \\ 0 & \lambda_2 & 0 & 0 & -\sqrt{\frac{3}{2}}\beta_2 \\ 0 & 0 & \lambda_3 & 0 & -\sqrt{\frac{3}{2}}\beta_3 \\ 0 & 0 & 0 & \lambda_4 & -\sqrt{\frac{3}{2}}\beta_4 \\ \sqrt{\frac{3}{2}}\beta_5 & \sqrt{\frac{3}{2}}\beta_6 & \sqrt{\frac{3}{2}}\beta_7 & \sqrt{\frac{3}{2}}\beta_8 & 0 \end{vmatrix}. \quad (12)$$

The matrices of the bilinear form should have the structure

$$\eta = \begin{vmatrix} \eta^{(1/2)} \otimes \gamma_4 & 0 \\ 0 & \eta^{(3/2)} \otimes I_2 \otimes \gamma_4 \end{vmatrix}, \quad \eta^{(1/2)} = \begin{vmatrix} k_1 & 0 & 0 & 0 & 0 \\ 0 & k_2 & 0 & 0 & 0 \\ 0 & 0 & k_3 & 0 & 0 \\ 0 & 0 & 0 & k_4 & 0 \\ 0 & 0 & 0 & 0 & k_5 \end{vmatrix}, \quad \eta^{(3/2)} = -\text{diag}(k_i), \quad (13)$$

where $k_i = \pm 1$. Due to the known constraint $(\eta\Gamma_4)^+ = \eta\Gamma_4$, we find restrictions on parameters λ_i, β_i :

$$\begin{aligned} \lambda_1^* &= \lambda_1, & \lambda_2^* &= \lambda_2, & \lambda_3^* &= \lambda_3, & \lambda_4^* &= \lambda_4, \\ \beta_5 &= -k_1 k_5 \beta_1^*, & \beta_6 &= -k_2 k_5 \beta_2^*, \\ \beta_7 &= -k_3 k_5 \beta_3^*, & \beta_8 &= -k_4 k_5 \beta_4^*. \end{aligned} \quad (14)$$

Because the particle under consideration has only one mass parameter M , and the spin equals $1/2$, the matrix $c^{(1/2)}$ (10) should have only one non-vanishing eigenvalue $(+1)$. We use this property to derive additional restrictions on parameters λ_i, β_i :

$$S_p c^{(1/2)} = 1, \quad S_p [(c^{(1/2)})^2] = 1, \quad S_p (c^{(1/2)})^3 = 1, \quad S_p (c^{(1/2)})^4 = 1, \quad \det c^{(1/2)} = 0. \quad (15)$$

From expression for $c^{(1/2)}$ (11) and condition (14), it follows

$$\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4 = 1. \quad (16)$$

From equality

$$(c^{(1/2)})^2 = \begin{vmatrix} \lambda_1^2 - \frac{3}{2}\beta_1\beta_5 & -\frac{3}{2}\beta_1\beta_6 & -\frac{3}{2}\beta_1\beta_7 & -\frac{3}{2}\beta_1\beta_8 & -\sqrt{\frac{3}{2}}\lambda_1\beta_1 \\ -\frac{3}{2}\beta_2\beta_5 & \lambda_2^2 - \frac{3}{2}\beta_2\beta_6 & -\frac{3}{2}\beta_2\beta_7 & -\frac{3}{2}\beta_2\beta_8 & -\sqrt{\frac{3}{2}}\lambda_2\beta_2 \\ -\frac{3}{2}\beta_3\beta_5 & -\frac{3}{2}\beta_3\beta_6 & \lambda_3^2 - \frac{3}{2}\beta_3\beta_7 & -\frac{3}{2}\beta_3\beta_8 & -\sqrt{\frac{3}{2}}\lambda_3\beta_3 \\ -\frac{3}{2}\beta_4\beta_5 & -\frac{3}{2}\beta_4\beta_6 & -\frac{3}{2}\beta_4\beta_7 & \lambda_4^2 - \frac{3}{2}\beta_4\beta_8 & -\sqrt{\frac{3}{2}}\lambda_4\beta_4 \\ \sqrt{\frac{3}{2}}\lambda_1\beta_5 & \sqrt{\frac{3}{2}}\lambda_2\beta_6 & \sqrt{\frac{3}{2}}\lambda_3\beta_7 & \sqrt{\frac{3}{2}}\lambda_4\beta_8 & -\frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) \end{vmatrix} \quad (17)$$

it follows (see (14))

$$(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \lambda_4^2) - 3(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) = 1. \quad (18)$$

Due to (15), we have $(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)^2 = 1$, or

$$(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \lambda_4^2) + 2(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) = 1. \quad (19)$$

From (17),(19) we obtain

$$(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) = 0. \quad (20)$$

From definition for $(C^{(1/2)})^3$:

$$(c^{(1/2)})^3 = \begin{vmatrix} \lambda_1(\lambda_1^2 - 3\beta_1\beta_5) \\ -\frac{3}{2}\beta_2\beta_5(\lambda_1 + \lambda_2) \\ -\frac{3}{2}\beta_3\beta_5(\lambda_1 + \lambda_3) \\ -\frac{3}{2}\beta_4\beta_5(\lambda_1 + \lambda_4) \\ \sqrt{\frac{3}{2}}\beta_5[\lambda_1^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \end{vmatrix}$$

$$\begin{aligned} & -\frac{3}{2}\beta_1\beta_6(\lambda_1 + \lambda_2) \\ & \lambda_2(\lambda_2^2 - 3\beta_2\beta_6) \\ & -\frac{3}{2}\beta_3\beta_6(\lambda_2 + \lambda_3) \\ & -\frac{3}{2}\beta_4\beta_6(\lambda_2 + \lambda_4) \\ & \sqrt{\frac{3}{2}}\beta_6[\lambda_2^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \end{aligned}$$

$$\begin{aligned} & -\frac{3}{2}\beta_1\beta_7(\lambda_1 + \lambda_3) \\ & -\frac{3}{2}\beta_2\beta_7(\lambda_2 + \lambda_3) \\ & \lambda_3(\lambda_3^2 - 3\beta_3\beta_7) \\ & -\frac{3}{2}\beta_4\beta_7(\lambda_3 + \lambda_4) \\ & \sqrt{\frac{3}{2}}\beta_7[\lambda_3^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \end{aligned}$$

$$\begin{aligned} & -\frac{3}{2}\beta_1\beta_8(\lambda_1 + \lambda_4) \\ & -\frac{3}{2}\beta_2\beta_8(\lambda_2 + \lambda_4) \\ & -\frac{3}{2}\beta_3\beta_8(\lambda_3 + \lambda_4) \\ & \lambda_4(\lambda_4^2 - 3\beta_4\beta_8) \\ & \sqrt{\frac{3}{2}}\beta_8[\lambda_4^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \end{aligned}$$

$$\begin{aligned} & -\sqrt{\frac{3}{2}}\beta_1[\lambda_1^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\ & -\sqrt{\frac{3}{2}}\beta_2[\lambda_2^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\ & -\sqrt{\frac{3}{2}}\beta_3[\lambda_3^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\ & -\sqrt{\frac{3}{2}}\beta_4[\lambda_4^2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\ & -\frac{3}{2}[\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8] \end{aligned} \quad (21)$$

we get

$$(\lambda_1^3 + \lambda_2^3 + \lambda_3^3) - \frac{9}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8) = 1. \quad (22)$$

Taking in mind (15), we find $(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)^3 = 1$, or

$$\begin{aligned} & (\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)^2 \\ & = (\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)[(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \lambda_4^2) + 2(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)] = 1, \end{aligned}$$

so we arrive at

$$\begin{aligned} & (\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) + 3\lambda_1(\lambda_2^2 + \lambda_3^2 + \lambda_4^2) + 3\lambda_2(\lambda_1^2 + \lambda_3^2 + \lambda_4^2) + 3\lambda_3(\lambda_1^2 + \lambda_2^2 + \lambda_4^2) \\ & + 3\lambda_4(\lambda_1^2 + \lambda_2^2 + \lambda_3^2) + 6(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) = 1. \end{aligned} \quad (23)$$

Relation (23) may be presented as follows

$$\begin{aligned} & (\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) - 3(\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) + 3(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \lambda_4^2) \\ & + 6(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) = 1, \end{aligned}$$

whence (taking in mind (16)) it follows

$$\begin{aligned} & -2(\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) + 3(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \lambda_4^2) \\ & + 6(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) = 1. \end{aligned} \quad (24)$$

Therefore, we get (taking in mind (19))

$$\begin{aligned} & (\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) = 3(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) \\ & - 3(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + 1. \end{aligned} \quad (25)$$

Allowing for (20), we transform (25) to other form

$$\begin{aligned} & (\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) = 3(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 \\ & + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) + \frac{9}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) + 1. \end{aligned} \quad (26)$$

From (26), (22) we derive

$$\begin{aligned}
 & (\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4) \\
 & + \frac{3}{2}[(1 - \lambda_1)\beta_1\beta_5 + (1 - \lambda_2)\beta_2\beta_6 + (1 - \lambda_3)\beta_3\beta_7 + (1 - \lambda_4)\beta_4\beta_8] = 0,
 \end{aligned} \tag{27}$$

or differently

$$\begin{aligned}
 & (\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4) \\
 & + \frac{3}{2}[(\lambda_2 + \lambda_3 + \lambda_4)\beta_1\beta_5 + (\lambda_1 + \lambda_3 + \lambda_4)\beta_2\beta_6 + (\lambda_1 + \lambda_2 + \lambda_4)\beta_3\beta_7 \\
 & + (\lambda_1 + \lambda_2 + \lambda_3)\beta_4\beta_8] = 0.
 \end{aligned} \tag{28}$$

Allowing for (25) and (22), we transform the last equation to the new form

$$\begin{aligned}
 & (\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4) - (\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4) \\
 & - \frac{3}{2}(\lambda_1 \beta_1 \beta_5 + \lambda_2 \beta_2 \beta_6 + \lambda_3 \beta_3 \beta_7 + \lambda_4 \beta_4 \beta_8) = 0.
 \end{aligned} \tag{29}$$

With the use of explicit form of $(C^{(1/2)})^4$ (we present the matrix by columns):

$$\begin{aligned}
 & (C^{(1/2)})^4 \\
 = & \begin{vmatrix}
 \lambda_1^4 - \frac{9}{2}\beta_1\beta_5[\lambda_1^2 - \frac{1}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_2\beta_5[\lambda_1^2 + \lambda_2^2 + \lambda_1\lambda_2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_3\beta_5[\lambda_1^2 + \lambda_3^2 + \lambda_1\lambda_3 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_4\beta_5[\lambda_1^2 + \lambda_4^2 + \lambda_1\lambda_4 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 \sqrt{\frac{3}{2}}\beta_5[\lambda_1^3 - \frac{3}{2}\lambda_1(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 \\
 -\frac{3}{2}\beta_1\beta_6[\lambda_1^2 + \lambda_2^2 + \lambda_1\lambda_2 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 \lambda_2^4 - \frac{9}{2}\beta_2\beta_6[\lambda_2^2 - \frac{1}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_3\beta_6[\lambda_2^2 + \lambda_3^2 + \lambda_2\lambda_3 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_4\beta_6[\lambda_2^2 + \lambda_4^2 + \lambda_2\lambda_4 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 \sqrt{\frac{3}{2}}\beta_6[\lambda_2^3 - \frac{3}{2}\lambda_2(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 \\
 -\frac{3}{2}\beta_1\beta_7[\lambda_1^2 + \lambda_3^2 + \lambda_1\lambda_3 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_2\beta_7[\lambda_2^2 + \lambda_3^2 + \lambda_2\lambda_3 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 \lambda_3^4 - \frac{9}{2}\beta_3\beta_7[\lambda_3^2 - \frac{1}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_4\beta_7[\lambda_3^2 + \lambda_4^2 + \lambda_3\lambda_4 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 \sqrt{\frac{3}{2}}\beta_7[\lambda_3^3 - \frac{3}{2}\lambda_3(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 \\
 -\frac{3}{2}\beta_1\beta_8[\lambda_1^2 + \lambda_4^2 + \lambda_1\lambda_4 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_2\beta_8[\lambda_2^2 + \lambda_4^2 + \lambda_2\lambda_4 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 -\frac{3}{2}\beta_3\beta_8[\lambda_3^2 + \lambda_4^2 + \lambda_3\lambda_4 - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 \lambda_4^4 - \frac{9}{2}\beta_4\beta_8[\lambda_4^2 - \frac{1}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)] \\
 \sqrt{\frac{3}{2}}\beta_8[\lambda_4^3 - \frac{3}{2}\lambda_4(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 \\
 -\sqrt{\frac{3}{2}}\beta_1[\lambda_1^3 - \frac{3}{2}\lambda_1(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 -\sqrt{\frac{3}{2}}\beta_2[\lambda_2^3 - \frac{3}{2}\lambda_2(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 -\sqrt{\frac{3}{2}}\beta_3[\lambda_3^3 - \frac{3}{2}\lambda_3(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 -\sqrt{\frac{3}{2}}\beta_4[\lambda_4^3 - \frac{3}{2}\lambda_4(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) - \frac{3}{2}(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8)] \\
 -\frac{3}{2}[(\lambda_1^2\beta_1\beta_5 + \lambda_2^2\beta_2\beta_6 + \lambda_3^2\beta_3\beta_7 + \lambda_4^2\beta_4\beta_8) - \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)^2]
 \end{vmatrix} \tag{30}
 \end{aligned}$$

we derive

$$\begin{aligned}
 & (\lambda_1^4 + \lambda_2^4 + \lambda_3^4 + \lambda_4^4) - 6(\lambda_1^2\beta_1\beta_5 + \lambda_2^2\beta_2\beta_6 + \lambda_3^2\beta_3\beta_7 + \lambda_4^2\beta_4\beta_8) \\
 & + \frac{9}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)^2 = 1.
 \end{aligned} \tag{31}$$

Taking in mind (20), the last equation (31) may be transformed to other form

$$(\lambda_1^4 + \lambda_2^4 + \lambda_3^4 + \lambda_4^4) - 6(\lambda_1^2\beta_1\beta_5 + \lambda_2^2\beta_2\beta_6 + \lambda_3^2\beta_3\beta_7 + \lambda_4^2\beta_4\beta_8) + 2(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)^2 = 1. \quad (32)$$

Whence, taking into account (16), we derive $(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)^4 = 1$, or differently

$$\begin{aligned} & (\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)^3 \\ &= (\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4) \left\{ (\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) + 3\lambda_1^2(\lambda_2 + \lambda_3 + \lambda_4) \right. \\ & \quad + 3\lambda_2^2(\lambda_1 + \lambda_3 + \lambda_4) + 3\lambda_3^2(\lambda_1 + \lambda_2 + \lambda_4) + 3\lambda_4^2(\lambda_2 + \lambda_3 + \lambda_4) \\ & \quad \left. + 6(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) \right\} \\ &= (\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4) \left\{ -2(\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) + 3\lambda_1^2(\lambda_2 + \lambda_3 + \lambda_4) \right. \\ & \quad \left. + 3(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \lambda_4^2) + 6(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) \right\} \\ &= -2(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)(\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) + 3\lambda_1^2(\lambda_2 + \lambda_3 + \lambda_4) \\ & \quad + 3(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)^2(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 + \lambda_4^2) + 6(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) \\ &= (\lambda_1^4 + \lambda_2^4 + \lambda_3^4 + \lambda_4^4) - 2\lambda_1(\lambda_2^3 + \lambda_3^3 + \lambda_4^3) - 2\lambda_2(\lambda_1^3 + \lambda_3^3 + \lambda_4^3) \\ & \quad - 2\lambda_3(\lambda_1^3 + \lambda_2^3 + \lambda_4^3) - 2\lambda_4(\lambda_1^3 + \lambda_2^3 + \lambda_3^3) + 6(\lambda_1^2\lambda_2^2 + \lambda_1^2\lambda_3^2 + \lambda_1^2\lambda_4^2 \\ & \quad + \lambda_2^2\lambda_3^2 + \lambda_2^2\lambda_4^2 + \lambda_3^2\lambda_4^2) + 6\lambda_1^3(\lambda_2 + \lambda_3 + \lambda_4) + 6\lambda_2^3(\lambda_1 + \lambda_3 + \lambda_4) \\ & \quad + 6\lambda_3^3(\lambda_1 + \lambda_2 + \lambda_4) + 6\lambda_4^3(\lambda_1 + \lambda_2 + \lambda_3) + 12\lambda_1^2(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) \\ & \quad + 12\lambda_2^2(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) + 12\lambda_3^2(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + 12\lambda_4^2(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3) + 24\lambda_1\lambda_2\lambda_3\lambda_4 \\ &= (\lambda_1^4 + \lambda_2^4 + \lambda_3^4 + \lambda_4^4) - 2\lambda_1(\lambda_2^3 + \lambda_3^3 + \lambda_4^3) - 2\lambda_2(\lambda_1^3 + \lambda_3^3 + \lambda_4^3) \\ & \quad - 2\lambda_3(\lambda_1^3 + \lambda_2^3 + \lambda_4^3) - 2\lambda_4(\lambda_1^3 + \lambda_2^3 + \lambda_3^3) + 6(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 \\ & \quad + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)^2 + 6\lambda_1^3(\lambda_2 + \lambda_3 + \lambda_4) + 6\lambda_2^3(\lambda_1 + \lambda_3 + \lambda_4) \\ & \quad + 6\lambda_3^3(\lambda_1 + \lambda_2 + \lambda_4) + 6\lambda_4^3(\lambda_1 + \lambda_2 + \lambda_3) \\ &= -3(\lambda_1^4 + \lambda_2^4 + \lambda_3^4 + \lambda_4^4) + 4(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)(\lambda_1^3 + \lambda_2^3 + \lambda_3^3 + \lambda_4^3) \\ & \quad + 6(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)^2 - 12\lambda_1\lambda_2\lambda_3\lambda_4 = 1. \end{aligned}$$

From the last equation, with the us of (15) and (25), we obtain

$$\begin{aligned} & (\lambda_1^4 + \lambda_2^4 + \lambda_3^4 + \lambda_4^4) = 1 + 4(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) \\ & \quad - 4(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) \\ & \quad + 2(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)^2 - 4\lambda_1\lambda_2\lambda_3\lambda_4, \end{aligned} \quad (33)$$

whence it follows (see (32), (33))

$$\begin{aligned} & (\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) + (\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)^2 \\ & \quad - (\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) \\ & \quad - \frac{3}{2}(\lambda_1^2\beta_1\beta_5 + \lambda_2^2\beta_2\beta_6 + \lambda_3^2\beta_3\beta_7 + \lambda_4^2\beta_4\beta_8) - 4\lambda_1\lambda_2\lambda_3\lambda_4 = 0. \end{aligned} \quad (34)$$

Equations (33) and (34) may be transformed to other form with the use of (20), (29):

$$\begin{aligned} & (\lambda_1^4 + \lambda_2^4 + \lambda_3^4 + \lambda_4^4) = 1 + 6(\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8) \\ & \quad + \frac{9}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)^2 - 4\lambda_1\lambda_2\lambda_3\lambda_4, \end{aligned} \quad (35)$$

$$\begin{aligned} & (\lambda_1\beta_1\beta_5 + \lambda_2\beta_2\beta_6 + \lambda_3\beta_3\beta_7 + \lambda_4\beta_4\beta_8) + \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)^2 \\ & \quad - (\lambda_1^2\beta_1\beta_5 + \lambda_2^2\beta_2\beta_6 + \lambda_3^2\beta_3\beta_7 + \lambda_4^2\beta_4\beta_8) - \frac{8}{3}\lambda_1\lambda_2\lambda_3\lambda_4 = 0. \end{aligned}$$

Due to the structure of eigenvalues of the matrix $C^{(1/2)}$ (10), its determinant equals zero:

$$\lambda_1\lambda_2(\lambda_4\beta_3\beta_7 + \lambda_3\beta_4\beta_8) + \lambda_3\lambda_4(\lambda_1\beta_2\beta_6 + \lambda_2\beta_1\beta_5) = 0. \quad (36)$$

The above restrictions on parameters λ_i, β_i may be found differently: by using the characteristic equation for the matrix $C^{(1/2)}$

$$(C^{(1/2)})^5 - a_1(C^{(1/2)})^4 + a_2(C^{(1/2)})^3 - a_3(C^{(1/2)})^2 + a_4(C^{(1/2)}) - a_5 = 0, \quad (37)$$

where

$$\begin{aligned} a_1 &= S_p(C^{(1/2)}), \quad a_2 = \frac{1}{2}[(S_p(C^{(1/2)}))^2 - S_p(C^{(1/2)})^2], \\ a_3 &= \frac{1}{3}\left\{S_p(C^{(1/2)})^3 - \frac{3}{2}(S_p(C^{(1/2)}))(S_p(C^{(1/2)})^2) + \frac{1}{2}(S_p C^{(1/2)})^3\right\}, \\ a_4 &= \frac{1}{4}\left\{S_p(C^{(1/2)})^4 - \frac{4}{3}(S_p C^{(1/2)})(S_p(C^{(1/2)})^3) \right. \\ &\quad \left. - \frac{1}{2}(S_p(C^{(1/2)})^2)^2 + (S_p C^{(1/2)})^2 S_p(C^{(1/2)})^2 - \frac{1}{6}(S_p C^{(1/2)})^4\right\}, \quad a_5 = \det C^{(1/2)}. \end{aligned} \quad (38)$$

Taking into account possible eigenvalues for matrix $C^{(1/2)}$, we get:

$$a_1 = 1, \quad a_2 = a_3 = a_4 = a_5 = 0. \quad (39)$$

At this we again arrive at relations (15),(20),(27),(37), and

$$\begin{aligned} 2\lambda_1\lambda_2\lambda_3\lambda_4 + 3\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + 3\beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) \\ + 3\beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + 3\beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3) = 0. \end{aligned} \quad (40)$$

2. Transformation to the spinor basis

Let us find explicit form of the matrix Γ_4 in spinor basis, this will permits us to determine an explicit form of the complete equation in the spinor basis. Such a transition is performed in several steps.

In canonical basis, the complete function Φ consists of the following components

$$\begin{aligned} \Psi^{(\text{канон.})} &= \left\{ (\Psi^{(1)})_{(0,1/2)}^{(0,1/2)}, (\Psi^{(1)})_{(0,-1/2)}^{(0,1/2)}, (\Psi^{(1)})_{(1/2,0)}^{(1/2,0)}, (\Psi^{(1)})_{(-1/2,0)}^{(1/2,0)}; \right. \\ &\quad (\Psi^{(2)})_{(0,1/2)}^{(0,1/2)}, (\Psi^{(2)})_{(0,-1/2)}^{(0,1/2)}, (\Psi^{(2)})_{(1/2,0)}^{(1/2,0)}, (\Psi^{(2)})_{(-1/2,0)}^{(1/2,0)}; \\ &\quad (\Psi^{(3)})_{(0,1/2)}^{(0,1/2)}, (\Psi^{(3)})_{(0,-1/2)}^{(0,1/2)}, (\Psi^{(3)})_{(1/2,0)}^{(1/2,0)}, (\Psi^{(3)})_{(-1/2,0)}^{(1/2,0)}; \\ &\quad (\Psi^{(4)})_{(0,1/2)}^{(0,1/2)}, (\Psi^{(4)})_{(0,-1/2)}^{(0,1/2)}, (\Psi^{(4)})_{(1/2,0)}^{(1/2,0)}, (\Psi^{(4)})_{(-1/2,0)}^{(1/2,0)}; \\ &\quad \Psi_{(1,1/2)}^{(1,1/2)}, \Psi_{(0,1/2)}^{(1,1/2)}, \Psi_{-1,1/2}^{(1,1/2)}, \Psi_{1,-1/2}^{(1,1/2)}, \Psi_{(0,-1/2)}^{(1,1/2)}, \Psi_{(-1,-1/2)}^{(1,1/2)}; \\ &\quad \left. \Psi_{1/2,1}^{(1/2,1)}, \Psi_{1/2,0}^{(1/2,1)}, \Psi_{1/2,-1}^{(1/2,1)}, \Psi_{-1/2,1}^{(1/2,1)}, \Psi_{-1/2,0}^{(1/2,1)}, \Psi_{-1/2,-1}^{(1/2,1)} \right\}_T. \end{aligned} \quad (41)$$

The functions $\Psi^{\text{G-Y(mod)}}$ in modified Gelfand – Yaglom basis and in canonical basis $\Psi^{(\text{can})}$ are related by the linear transformation A:

$$\Psi^{\text{G-Y(mod)}} = A\Psi^{(\text{can})}. \quad (42)$$

Explicate form of the matrix A according to the general rule

$$\Psi_{s,s_3}^{(l,l')} = \sum_{l_3,l'_3} C_{l_3,l'_3}^{s,s_3} \Psi_{(l_3,l'_3)}^{(l,l')}, \quad (43)$$

where C_{l_3,l'_3}^{s,s_3} are the Clebsch – Gordan coefficients; the summing is over indices l_3, l'_3 at restriction $l_3 + l'_3 = s_3$. For given case, we obtain

$$A = \begin{vmatrix} I_{16 \times 16} & & & & & & & & & & & & & & & & \\ & -\frac{1}{\sqrt{3}} & & \frac{2}{\sqrt{3}} & & & & & & & & & & & & & \\ & & -\sqrt{\frac{2}{3}} & & \frac{1}{\sqrt{3}} & & & & & & & & & & & & \\ & & & & & & \frac{1}{\sqrt{3}} & & -\sqrt{\frac{2}{3}} & & & & & & & & \\ & 1 & & & & & & & \sqrt{\frac{2}{3}} & & -\frac{1}{\sqrt{3}} & & & & & & \\ & & & & & 1 & & & & & & & & & & & \\ & & & & & & 1 & & & & & & & & & & \\ & & & & & & & & & & & & & & & 1 & \\ & & \sqrt{\frac{2}{3}} & & \frac{1}{\sqrt{3}} & & & & & & & & & & & & \\ & & & \frac{1}{\sqrt{3}} & & \sqrt{\frac{2}{3}} & & & & & & & & & & & \\ & & & & & & & \sqrt{\frac{2}{3}} & & \frac{1}{\sqrt{3}} & & & & & & & \\ & & & & & & & & \frac{1}{\sqrt{3}} & & \sqrt{\frac{2}{3}} & & & & & & \\ & & & & & & & & & \frac{1}{\sqrt{3}} & & \sqrt{\frac{2}{3}} & & & & & \end{vmatrix},$$

$$A^+ = \begin{vmatrix} I_{16 \times 16} & & & & & & & & & & & & & & & & \\ & & & & & 1 & & & & & & & & & & & \\ & -\frac{1}{\sqrt{3}} & & & & & & & \sqrt{\frac{2}{3}} & & & & & & & & \\ & & -\sqrt{\frac{2}{3}} & & & & & & & \frac{1}{\sqrt{3}} & & & & & & & \\ & \sqrt{\frac{2}{3}} & & & & & & & \frac{1}{\sqrt{3}} & & & & & & & & \\ & & \frac{1}{\sqrt{3}} & & & & & & & \sqrt{\frac{2}{3}} & & & & & & & \\ & & & & & 1 & & & & & & & & & & & \\ & & & & & & 1 & & & & & & & & & & \\ & & & & \frac{1}{\sqrt{3}} & & & & & & & & \sqrt{\frac{2}{3}} & & & & \\ & & & \sqrt{\frac{2}{3}} & & & & & & & & & & \frac{1}{\sqrt{3}} & & & \\ & & & & -\sqrt{\frac{2}{3}} & & & & & & & & & \frac{1}{\sqrt{3}} & & & \\ & & & & & -\frac{1}{\sqrt{3}} & & & & & & & & & \sqrt{\frac{2}{3}} & & \\ & & & & & & 1 & & & & & & & & & & \end{vmatrix},$$

and

$$\Psi^{(\text{can})} = A^+ \Psi^{\text{G-Y(mod)}}, \quad \Gamma_4^{(\text{can})} = A^+ \Gamma_4^{\text{G-Y(mod)}} A.$$

In detailed form, the matrix $\Gamma_4^{(\text{G-Y(mod)})}$ reads

$$\Gamma_4^{(\text{spin})} = B^{-1} \Gamma_4^{(\text{can})} B = \begin{vmatrix} \Gamma_{16 \times 16} & \Gamma_{12 \times 16} \\ \Gamma_{16 \times 12} & 0 \end{vmatrix}, \quad (44)$$

$$\Gamma_{16 \times 16} = \begin{vmatrix} & & \lambda_1 & & & & & & & & & & & & & \\ & & & \lambda_1 & & & & & & & & & & & & \\ \lambda_1 & & & & & & & & & & & & & & & \\ & \lambda_1 & & & & & & & & & & & & & & \\ & & & & & \lambda_2 & & & & & & & & & & \\ & & & & & & \lambda_2 & & & & & & & & & \\ & & & & \lambda_2 & & & & & & & & & & & \\ & & & & & \lambda_2 & & & & & & & & & & \\ & & & & & & & \lambda_3 & & & & & & & & \\ & & & & & & & & \lambda_3 & & & & & & & \\ & & & & & & \lambda_3 & & & & & & & & & \\ & & & & & & & \lambda_3 & & & & & & & & \\ & & & & & & & & & \lambda_4 & & & & & & \\ & & & & & & & & & & \lambda_4 & & & & & \\ & & & & & & & & & & & \lambda_4 & & & & \\ & & & & & & & & & & & & \lambda_4 & & & \end{vmatrix},$$

[illegible]

therefore

$$\Gamma_4^{(\text{can})} = \begin{vmatrix} \Gamma_{16 \times 16} & \Gamma_{12 \times 16} \\ \Gamma_{16 \times 12} & 0 \end{vmatrix} = A^+ \Gamma_4^{\text{G-Y(mod)}} A, \quad (45)$$

where

[illegible]

$$\Gamma_{12 \times 16} = \begin{vmatrix} \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\frac{\beta_1}{\sqrt{2}} & \cdot & \beta_1 & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\beta_1 & \cdot & \frac{\beta_1}{\sqrt{2}} & \cdot \\ \cdot & \frac{\beta_1}{\sqrt{2}} & \cdot & -\beta_1 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \beta_1 & \cdot & -\frac{\beta_1}{\sqrt{2}} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\frac{\beta_2}{\sqrt{2}} & \cdot & \beta_2 & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\beta_2 & \cdot & \frac{\beta_2}{\sqrt{2}} & \cdot & \cdot \\ \cdot & \frac{\beta_2}{\sqrt{2}} & \cdot & -\beta_2 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \beta_2 & \cdot & -\frac{\beta_2}{\sqrt{2}} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\frac{\beta_3}{\sqrt{2}} & \cdot & \beta_3 & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\beta_3 & \cdot & \frac{\beta_3}{\sqrt{2}} & \cdot & \cdot \\ \cdot & \frac{\beta_3}{\sqrt{2}} & \cdot & -\beta_3 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \beta_3 & \cdot & -\frac{\beta_3}{\sqrt{2}} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\frac{\beta_4}{\sqrt{2}} & \cdot & \beta_4 & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & -\beta_4 & \cdot & \frac{\beta_4}{\sqrt{2}} & \cdot & \cdot \\ \cdot & \frac{\beta_4}{\sqrt{2}} & \cdot & -\beta_4 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \beta_4 & \cdot & -\frac{\beta_4}{\sqrt{2}} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \end{vmatrix},$$

$$\Gamma_{16 \times 12} = \begin{vmatrix} \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & -\frac{\beta_5}{\sqrt{2}} & \cdot & \cdot & \cdot & -\frac{\beta_6}{\sqrt{2}} & \cdot & \cdot & -\frac{\beta_7}{\sqrt{2}} & \cdot & -\frac{\beta_8}{\sqrt{2}} \\ \cdot & \cdot & \cdot & -\beta_5 & \cdot & \cdot & -\beta_6 & \cdot & \cdot & -\beta_7 & \cdot & -\beta_8 \\ \cdot & \cdot & \beta_5 & \cdot & \cdot & \beta_6 & \cdot & \cdot & \beta_7 & \cdot & \cdot & \beta_8 \\ \cdot & \cdot & \cdot & \frac{\beta_5}{\sqrt{2}} & \cdot & \cdot & \frac{\beta_6}{\sqrt{2}} & \cdot & \cdot & \frac{\beta_7}{\sqrt{2}} & \cdot & \frac{\beta_8}{\sqrt{2}} \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \frac{\beta_5}{\sqrt{2}} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \beta_5 & \cdot & \cdot & \beta_6 & \cdot & \cdot & \beta_7 & \cdot & \cdot & \beta_8 & \cdot \\ -\beta_5 & \cdot & \cdot & -\beta_6 & \cdot & \cdot & -\beta_7 & \cdot & \cdot & -\beta_8 & \cdot & \cdot \\ \cdot & -\frac{\beta_5}{\sqrt{2}} & \cdot & \cdot & -\frac{\beta_6}{\sqrt{2}} & \cdot & \cdot & -\frac{\beta_7}{\sqrt{2}} & \cdot & \cdot & -\frac{\beta_8}{\sqrt{2}} & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \end{vmatrix}.$$

When defining the matrix $\Gamma_4^{(\text{spin})}$ in spinor basis, we apply the general formula

$$\Psi_{(l_3, l'_3)}^{(l, l')} = \left[\frac{(2l)!}{(l+l_3)!(l-l_3)!} \right]^{1/2} \left[\frac{(2l')!}{(l'+l'_3)!(l'-l'_3)!} \right]^{1/2} \Psi_{(1\dots 12\dots 2)}^{(i\dots i2\dots 2)}. \quad (46)$$

As the result we obtain the following linear transformation

$$\Psi^{(\text{can})} = B \Psi^{(\text{spin})}, \quad (47)$$

where

$$\begin{aligned} \Psi^{(\text{spin})} = & \left\{ (\Psi^{(1)})^i, (\Psi^{(1)})^{\dot{2}}, (\Psi^{(1)})_1, (\Psi^{(1)})_2; (\Psi^{(2)})^i, (\Psi^{(2)})^{\dot{2}}, (\Psi^{(2)})_1, \right. \\ & (\Psi^{(2)})_2; (\Psi^{(3)})^i, (\Psi^{(3)})^{\dot{2}}, (\Psi^{(3)})_1, (\Psi^{(3)})_2; (\Psi^{(4)})^i, (\Psi^{(4)})^{\dot{2}}, \\ & (\Psi^{(4)})_1, (\Psi^{(4)})_2; \Psi_{(11)}^i, \Psi_{(12)}^i, \Psi_{(22)}^i; \Psi_{(11)}^{\dot{2}}, \Psi_{(12)}^{\dot{2}}, \Psi_{(22)}^{\dot{2}}; \\ & \left. \Psi_1^{(ii)}, \Psi_1^{(i\dot{2})}, \Psi_1^{(\dot{2}\dot{2})}, \Psi_2^{(ii)}, \Psi_2^{(i\dot{2})}, \Psi_2^{(\dot{2}\dot{2})} \right\}, \end{aligned} \quad (48)$$

$$B = \begin{vmatrix} I_{16 \times 16} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \sqrt{2} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \sqrt{2} & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \sqrt{2} & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot & 1 \end{vmatrix},$$

The matrix $\Gamma_4^{(\text{spin})}$ may be presented with the use of the spinor elements of the complete matrix algebra:

$$\begin{aligned} \Gamma_4^{\text{spin}} = \frac{1}{i} \Big\{ & \lambda_1 \left[(\sigma^4)^{\dot{a}b} e_{\dot{a}1}^0 \otimes e_{b1}^0 + (\sigma^4)_{ab} e_{\dot{0}1}^1 \otimes e_{01}^1 \right] \lambda_2 \left[(\sigma^4)^{\dot{a}b} e_{\dot{a}2}^0 \otimes e_{b2}^0 + (\sigma^4)_{ab} e_{\dot{0}2}^1 \otimes e_{02}^1 \right] \\ & + \lambda_3 \left[(\sigma^4)^{\dot{a}b} e_{\dot{a}3}^0 \otimes e_{b3}^0 + (\sigma^4)_{ab} e_{\dot{0}3}^1 \otimes e_{03}^1 \right] \lambda_4 \left[(\sigma^4)^{\dot{a}b} e_{\dot{a}4}^0 \otimes e_{b4}^0 + (\sigma^4)_{ab} e_{\dot{0}4}^1 \otimes e_{04}^1 \right] \\ & + \beta_1 (\sigma^4)_{\dot{c}}^b [e_{\dot{a}1}^{\dot{c}} \otimes e_b^0 + e_{\dot{0}1}^{\dot{c}} \otimes e_{(ba)}^1] \beta_2 (\sigma^4)_{\dot{c}}^b [e_{\dot{a}2}^{\dot{c}} \otimes e_b^0 + e_{\dot{0}2}^{\dot{c}} \otimes e_{(ba)}^1] \\ & + \beta_3 (\sigma^4)_{\dot{c}}^b [e_{\dot{a}3}^{\dot{c}} \otimes e_b^0 + e_{\dot{0}3}^{\dot{c}} \otimes e_{(ba)}^1] + \beta_4 (\sigma^4)_{\dot{c}}^b [e_{\dot{a}4}^{\dot{c}} \otimes e_b^0 + e_{\dot{0}4}^{\dot{c}} \otimes e_{(ba)}^1] + \beta_5 (\sigma^4)_a^{\dot{d}} [e_{(\dot{a}b)}^{\dot{d}} \otimes e_{01}^a + e_{\dot{d}}^{\dot{0}1} \otimes e_{c1}^{(ac)}] \\ & + \beta_6 (\sigma^4)_a^{\dot{d}} [e_{(\dot{a}b)}^{\dot{d}} \otimes e_{02}^a + e_{\dot{d}}^{\dot{0}2} \otimes e_{c2}^{(ac)}] + \beta_7 (\sigma^4)_a^{\dot{d}} [e_{(\dot{a}b)}^{\dot{d}} \otimes e_{03}^a + e_{\dot{d}}^{\dot{0}3} \otimes e_{c3}^{(ac)}] \\ & + \beta_8 (\sigma^4)_a^{\dot{d}} [e_{(\dot{a}b)}^{\dot{d}} \otimes e_{04}^a + e_{\dot{d}}^{\dot{0}4} \otimes e_{c4}^{(ac)}] \Big\}. \end{aligned} \quad (50)$$

Similarly may be obtained all four matrices $\Gamma_\mu^{(\text{spin})}$; to this end, it suffices to make the formal change $(\sigma^4) \Rightarrow (\sigma^\mu)$:

$$(\sigma^1)_{\dot{a}b} = \begin{vmatrix} 0 & 1 \\ 1 & 0 \end{vmatrix}_{\dot{a}b}, \quad (\sigma^2)_{\dot{a}b} = \begin{vmatrix} 0 & -i \\ i & 0 \end{vmatrix}_{\dot{a}b}, \quad (\sigma^3)_{\dot{a}b} = \begin{vmatrix} 1 & 0 \\ 0 & -1 \end{vmatrix}_{\dot{a}b}, \quad (\sigma^4)_{\dot{a}b} = \begin{vmatrix} i & 0 \\ 0 & i \end{vmatrix}_{\dot{a}b}.$$

Recall that the spinor elements of the complete matrix algebra obey the following rules

$$(e_B^A)_{\dot{D}}^{\dot{C}} = \delta_D^A \delta_{\dot{D}}^{\dot{C}}, \quad e_B^A e_D^{\dot{C}} = \delta_D^A e_B^{\dot{C}}, \quad (e_B^A)^C_D = \delta_D^A \delta_B^C, \quad e_B^A e_D^C = \delta_B^C e_D^A, \quad (51)$$

the digital index at spinors (like as $\dot{a}_1$) determine the number of spinor equation with respect to repeating presentations.

The complete matrix equation determined by four $\Gamma_\mu^{(\text{spin})}$ consists of 5 equations:

$$\begin{aligned} & \lambda_1 \left\{ \partial^{\dot{a}b} e_a^{\dot{0}} \otimes e_b^0 + \partial_{ab} e_{\dot{0}}^b \otimes e_0^a \right\} \Psi^{(1)} + \beta_1 \partial_c^b \left\{ e_a^{(\dot{a}c)} \otimes e_b^0 \right. \\ & \left. + e_{\dot{0}}^c \otimes e_{(ba)}^a \right\} \Psi + M \left\{ e_{\dot{0}}^0 \otimes e_a^a + e_{\dot{a}}^a \otimes e_0^0 \right\} \Psi^{(1)} = 0, \end{aligned} \quad (52)$$

$$\begin{aligned} & \lambda_2 \left\{ \partial^{\dot{a}b} e_a^{\dot{0}} \otimes e_b^0 + \partial_{ab} e_{\dot{0}}^b \otimes e_0^a \right\} \Psi^{(2)} + \beta_2 \partial_c^b \left\{ e_a^{(\dot{a}c)} \otimes e_b^0 + e_{\dot{0}}^c \otimes e_{(ba)}^a \right\} \Psi \\ & + M \left\{ e_{\dot{a}}^a \otimes e_0^0 + e_{\dot{0}}^0 \otimes e_a^a \right\} \Psi^{(2)} = 0, \end{aligned} \quad (53)$$

$$\begin{aligned} & \lambda_3 \left(\partial^{\dot{a}b} e_a^{\dot{0}} \otimes e_b^0 + \partial_{ab} e_{\dot{0}}^b \otimes e_0^a \right) \Psi^{(3)} + \beta_3 \partial_c^b \left(e_a^{(\dot{a}c)} \otimes e_b^0 + e_{\dot{0}}^c \otimes e_{(ba)}^a \right) \Psi \\ & + M \left\{ e_{\dot{a}}^a \otimes e_0^0 + e_{\dot{0}}^0 \otimes e_a^a \right\} \Psi^{(3)} = 0, \end{aligned} \quad (54)$$

$$\begin{aligned} & \lambda_4 \left(\partial^{\dot{a}b} e_a^{\dot{0}} \otimes e_b^0 + \partial_{ab} e_{\dot{0}}^b \otimes e_0^a \right) \Psi^{(4)} + \beta_4 \partial_c^b \left(e_a^{(\dot{a}c)} \otimes e_b^0 + e_{\dot{0}}^c \otimes e_{(ba)}^a \right) \Psi \\ & + M \left\{ e_{\dot{a}}^a \otimes e_0^0 + e_{\dot{0}}^0 \otimes e_a^a \right\} \Psi^{(4)} = 0, \end{aligned} \quad (55)$$

$$\begin{aligned} & \partial_a^{\dot{d}} \left(e_{(\dot{d}b)}^{\dot{b}} \otimes e_0^a + e_d^0 \otimes e_c^{(ac)} \right) \left(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} \right) \\ & + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)} + M \left(e_{(\dot{a}b)}^{(\dot{a}b)} \otimes e_c^c + e_{\dot{a}}^a \otimes e_{(bc)}^{(bc)} \right) \Psi = 0, \end{aligned} \quad (56)$$

where

$$\partial_{\dot{a}b} = \frac{1}{i} \partial_\mu (\sigma^\mu)_{\dot{a}b}. \quad (57)$$

Let us transform the above system (52)–(56) to spinor form. We start with Eq. (52), it leads to (see (51))

$$\begin{aligned} & \lambda_1 \left\{ \partial^{\dot{a}b} (e_a^{\dot{0}})^{\dot{A}}_B (e_b^0)^C_D \Psi^{(1)\dot{B}}_C + \partial_{ab} (e_{\dot{0}}^b)^{\dot{A}}_B (e_0^a)^C_D \Psi^{(1)\dot{B}}_D \right\} + \beta_1 \partial_c^b \left\{ (e_a^{(\dot{a}c)})^{\dot{A}}_B (e_b^0)^C_D \Psi^{\dot{B}}_C + (e_{\dot{0}}^c)^{\dot{A}}_B (e_{(ba)}^a)^C_D \Psi^{\dot{B}}_C \right\} \\ & + M \left\{ (e_{\dot{a}}^a)^{\dot{A}}_B (e_0^0)^C_D (\Psi^{(1)})^{\dot{B}}_C + (e_{\dot{0}}^0)^{\dot{A}}_B (e_a^a)^C_D (\Psi^{(1)})^{\dot{B}}_C \right\} = 0, \end{aligned}$$

or

$$\begin{aligned} \lambda_1 \left\{ \partial^{\dot{a}b} \delta_a^{\dot{A}} \delta_D^0 \Psi^{(1)\dot{b}}_b + \partial_{ab} \delta_0^{\dot{A}} \delta_D^a \Psi^{(1)\dot{b}}_0 \right\} + \beta_1 \partial_{\dot{c}}^b \left\{ \delta_a^{\dot{A}} \delta_D^0 \Psi^{(\dot{a}\dot{c})}_b + \delta_0^{\dot{A}} \delta_D^a \Psi^{(\dot{c})}_{(ba)} \right\} \\ + M \left\{ \delta_a^{\dot{A}} \delta_D^0 (\Psi^{(1)})_{\dot{a}}^0 + \delta_0^{\dot{A}} \delta_D^a \Psi^{(1)\dot{a}}_a \right\} = 0. \end{aligned}$$

Let $\dot{A} = \dot{k}$, $D = 0$, then we obtain

$$\lambda_1 \partial^{\dot{k}b} \Psi_b^{(1)} + \beta_1 \partial_{\dot{c}}^b \Psi_b^{(\dot{k}\dot{c})} + M \Psi^{(1)\dot{k}} = 0. \quad (58)$$

Let $\dot{A} = \dot{0}$, $D = k$, then we get

$$\lambda_1 \partial_{kb} \Psi^{(1)\dot{b}} + \beta_1 \partial_{\dot{c}}^b \Psi_{(bk)}^{\dot{c}} + M \Psi_k^{(1)} = 0. \quad (59)$$

Taking in account identities

$$\Psi_c^{(\dot{a}\dot{b})} = \delta_{(\dot{r}\dot{n})}^{(\dot{a}\dot{b})} (\sigma^\mu)_c^{\dot{r}} \Psi_\mu^{\dot{n}} = \frac{1}{2} (\delta_r^{\dot{a}} \delta_n^{\dot{b}} + \delta_n^{\dot{a}} \delta_r^{\dot{b}}) (\sigma^\mu)_c^{\dot{r}} \Psi_\mu^{\dot{n}} = \frac{1}{2} \left\{ (\sigma^\mu)_c^{\dot{a}} \Psi_\mu^{\dot{b}} + (\sigma^\mu)_c^{\dot{b}} \Psi_\mu^{\dot{a}} \right\}, \quad (60)$$

$$\Psi_{(ab)}^{\dot{c}} = \delta_{(ab)}^{(rn)} (\sigma^\mu)_r^{\dot{c}} \Psi_{n\mu} = \frac{1}{2} (\delta_a^r \delta_b^n + \delta_a^n \delta_b^r) (\sigma^\mu)_r^{\dot{c}} \Psi_{n\mu} = \frac{1}{2} \left\{ (\sigma^\mu)_a^{\dot{c}} \Psi_{b\mu} + (\sigma^\mu)_b^{\dot{c}} \Psi_{a\mu} \right\}, \quad (61)$$

we derive (see (58) and (59))

$$\lambda_1 \partial^{\dot{k}b} \Psi_b^{(1)} + \frac{\beta_1}{2} \left\{ (\sigma^\mu)_b^{\dot{k}} \partial_c^b \Psi_\mu^{\dot{c}} + \partial_{\dot{c}}^b (\sigma^\mu)_b^{\dot{c}} \Psi_\mu^{\dot{k}} \right\} + M \Psi^{(1)\dot{k}} = 0,$$

$$\lambda_1 \partial_{kb} \Psi^{(1)\dot{b}} + \frac{\beta_1}{2} \left\{ (\sigma^\mu)_k^{\dot{c}} \partial_{\dot{c}}^b \Psi_{b\mu} + \partial_{\dot{c}}^b (\sigma^\mu)_b^{\dot{c}} \Psi_{k\mu} \right\} + M \Psi_k^{(1)} = 0,$$

or

$$\lambda_1 \partial^{\dot{k}b} \Psi_b^{(1)} + \frac{\beta_1}{2} \left\{ -(\sigma^\mu)^{\dot{k}b} \partial_{b\dot{c}} \Psi_\mu^{\dot{c}} - \partial^{\dot{c}b} (\sigma^\mu)_{b\dot{c}} \Psi_\mu^{\dot{k}} \right\} + M \Psi^{(1)\dot{k}} = 0, \quad (62)$$

$$\lambda_1 \partial_{kb} \Psi^{(1)\dot{b}} + \frac{\beta_1}{2} \left\{ -(\sigma^\mu)_{k\dot{c}} \partial^{\dot{c}b} \Psi_{b\mu} - \partial_{b\dot{c}} (\sigma^\mu)^{\dot{c}b} \Psi_{k\mu} \right\} + M \Psi_k^{(1)} = 0. \quad (63)$$

Joining equations (62) and (63) into a matrix one one, we get

$$\begin{aligned} \lambda_1 \left| \begin{array}{cc} 0 & \partial^{\dot{k}b} \\ \partial_{kb} & 0 \end{array} \right| \left| \begin{array}{c} \Psi^{(1)\dot{b}} \\ \Psi_b^{(1)} \end{array} \right| + \frac{\beta_1}{2} \left\{ - \left| \begin{array}{cc} (\sigma^\mu)^{\dot{k}b} \partial_{b\dot{c}} & 0 \\ 0 & (\sigma^\mu)_{k\dot{c}} \partial^{\dot{c}b} \end{array} \right| \left| \begin{array}{c} \Psi_\mu^{\dot{c}} \\ \Psi_{c\mu} \end{array} \right| \right. \\ \left. - \left| \begin{array}{cc} \partial^{\dot{c}b} (\sigma^\mu)_{b\dot{c}} \delta_{\dot{c}}^{\dot{k}} & 0 \\ 0 & \partial_{a\dot{b}} (\sigma^\mu)^{\dot{b}a} \delta_{\dot{c}}^{\dot{k}} \end{array} \right| \left| \begin{array}{c} \Psi_\mu^{\dot{c}} \\ \Psi_{c\mu} \end{array} \right| \right\} + M \left| \begin{array}{c} \Psi^{(1)\dot{k}} \\ \Psi_k^{(1)} \end{array} \right| = 0. \end{aligned} \quad (64)$$

Allowing for the identities

$$\sigma^{\mu\dot{a}b} \sigma_{b\dot{c}}^\nu + \sigma^{\nu\dot{a}b} \sigma_{b\dot{c}}^\mu = -2\delta_{\mu\nu} \delta_{\dot{c}}^{\dot{a}}, \quad \sigma_{a\dot{b}}^\mu \sigma^{\nu\dot{b}c} + \sigma_{a\dot{b}}^\nu \sigma^{\mu\dot{b}c} = -2\delta_{\mu\nu} \delta_a^c, \quad \sigma^{\mu\dot{a}b} \sigma_{b\dot{a}}^\nu = \sigma_{a\dot{b}}^\mu \sigma^{\nu\dot{b}a} = -2\delta_{\mu\nu}, \quad (65)$$

we get the following presentation for Dirac matrices

$$\frac{1}{i} \left| \begin{array}{cc} 0 & \sigma^{\mu\dot{a}b} \\ \partial_{ab}^\mu & 0 \end{array} \right| = \gamma_\mu; \quad (66)$$

so we arrive at

$$\lambda_1 \hat{\partial} \Psi^{(1)} - 2\beta_1 \left\{ -\partial_\mu \Psi_\mu - \frac{1}{4} \hat{\partial} (\gamma_\mu \Psi_\mu) \right\} + M \Psi^{(1)} = 0, \quad (67)$$

where bispinors and vector-bispinor are introduced:

$$\Psi^{(1)} = \left| \begin{array}{c} \Psi^{(1)a'} \\ \Psi_a^{(1)} \end{array} \right|, \quad \Psi^{(1)} = \left| \begin{array}{c} \Psi_\mu^{\dot{a}} \\ \Psi_{a\mu} \end{array} \right|.$$

In similar way, we derive the following equations

$$\lambda_2 \partial^{kb} \Psi_b^{(2)} + \beta_2 \partial_c^b \Psi_b^{(k\dot{c})} + M \Psi^{(2)\dot{k}} = 0, \quad (68)$$

$$\lambda_2 \partial_{kb} \Psi^{(2)\dot{b}} + \beta_2 \partial_b^c \Psi_{(kc)}^{\dot{b}} + M \Psi_k^{(2)} = 0, \quad (69)$$

$$\lambda_2 \hat{\partial} \Psi^{(2)} - i2\beta_2 \left\{ \partial_\mu \Psi_\mu - \frac{1}{4} \hat{\partial}(\gamma_\mu \Psi_\mu) \right\} + M \Psi^{(2)} = 0; \quad (70)$$

$$\lambda_3 \partial^{kb} \Psi_b^{(3)} + \beta_3 \partial_c^b \Psi_b^{(k\dot{c})} + M \Psi^{(3)\dot{k}} = 0, \quad (71)$$

$$\lambda_3 \partial_{kb} \Psi^{(3)\dot{b}} + \beta_3 \partial_b^c \Psi_{(kc)}^{\dot{b}} + M \Psi_k^{(3)} = 0, \quad (72)$$

$$\lambda_3 \hat{\partial} \Psi^{(3)} - i2\beta_3 \left\{ \partial_\mu \Psi_\mu - \frac{1}{4} \hat{\partial}(\gamma_\mu \Psi_\mu) \right\} + M \Psi^{(3)} = 0; \quad (73)$$

$$\lambda_4 \partial^{kb} \Psi_b^{(4)} + \beta_4 \partial_c^b \Psi_b^{(k\dot{c})} + M \Psi^{(4)\dot{k}} = 0, \quad (74)$$

$$\lambda_4 \partial_{kb} \Psi^{(4)\dot{b}} + \beta_4 \partial_b^c \Psi_{(kc)}^{\dot{b}} + M \Psi_k^{(4)} = 0, \quad (75)$$

$$\lambda_4 \hat{\partial} \Psi^{(4)} - i2\beta_4 \left\{ \partial_\mu \Psi_\mu - \frac{1}{4} \hat{\partial}(\gamma_\mu \Psi_\mu) \right\} + M \Psi^{(4)} = 0. \quad (76)$$

Now we turn to Eq. (56), it leads to

$$\begin{aligned} & \partial_a^{\dot{d}} \left\{ (e_{(\dot{a}\dot{b})})_{\dot{B}}^{\dot{A}} (e_0^a)_D^C + (e_d^0)_{\dot{B}}^{\dot{A}} (e_c^{(ac)})_D^C \right\} \left\{ \beta_5 \Psi^{(1)\dot{B}}_C \right. \\ & \left. + \beta_6 \Psi^{(2)\dot{B}}_C + \beta_7 \Psi^{(3)\dot{B}}_C + \beta_8 \Psi^{(4)\dot{B}}_C \right\} + M \left\{ (e_{(\dot{a}\dot{b})})_D^C (e_c^{\dot{a}})_{\dot{B}}^{\dot{A}} (e_{(bc)})_D^C \right\} \Psi_C^{\dot{B}} = 0, \end{aligned}$$

or

$$\begin{aligned} & \partial_a^{\dot{d}} \delta_{(\dot{d}\dot{b})}^{\dot{A}} \delta_D^a (\beta_5 \Psi_0^{(1)\dot{b}} + \beta_6 \Psi_0^{(2)\dot{b}} + \beta_7 \Psi_0^{(3)\dot{b}} + \beta_8 \Psi_0^{(4)\dot{b}}) \\ & + \partial_a^{\dot{d}} \delta_{\dot{d}}^{\dot{A}} \delta_D^{(ac)} (\beta_5 \Psi_c^{(1)\dot{0}} + \beta_6 \Psi_c^{(2)\dot{0}} + \beta_7 \Psi_c^{(3)\dot{0}} + \beta_8 \Psi_c^{(4)\dot{0}}) + M \left\{ \delta_{(\dot{a}\dot{b})}^{\dot{A}} \delta_D^c \Psi_c^{(\dot{a}\dot{b})} + \delta_{\dot{a}}^{\dot{A}} \delta_D^{(bc)} \Psi_{(bc)}^{\dot{a}} \right\} = 0. \end{aligned} \quad (77)$$

Let $\dot{A} = (\dot{k}\dot{n})$, $D = r$; then we get

$$\partial_r^{\dot{d}} \delta_{(\dot{d}\dot{b})}^{\dot{A}} (\beta_5 \Psi^{(1)\dot{b}} + \beta_6 \Psi^{(2)\dot{b}} + \beta_7 \Psi^{(3)\dot{b}} + \beta_8 \Psi^{(4)\dot{b}} + M \Psi_r^{(\dot{k}\dot{n})}) = 0,$$

or

$$\frac{1}{2} \partial_r^{\dot{d}} (\delta_{\dot{d}}^{\dot{k}} \delta_{\dot{b}}^{\dot{n}} + \delta_{\dot{b}}^{\dot{k}} \delta_{\dot{d}}^{\dot{n}}) (\beta_5 \Psi^{(1)\dot{b}} + \beta_6 \Psi^{(2)\dot{b}} + \beta_7 \Psi^{(3)\dot{b}} + \beta_8 \Psi^{(4)\dot{b}}) + M \Psi_r^{(\dot{k}\dot{n})} = 0,$$

that is

$$\begin{aligned} & \frac{1}{2} \left\{ \partial_r^{\dot{k}} (\beta_5 \Psi^{(1)\dot{n}} + \beta_6 \Psi^{(2)\dot{n}} + \beta_7 \Psi^{(3)\dot{n}} + \beta_8 \Psi^{(4)\dot{n}}) \right. \\ & \left. + \partial_r^{\dot{n}} (\beta_5 \Psi^{(1)\dot{k}} + \beta_6 \Psi^{(2)\dot{k}} + \beta_7 \Psi^{(3)\dot{k}} + \beta_8 \Psi^{(4)\dot{k}}) \right\} + M \Psi_r^{(\dot{k}\dot{n})} = 0. \end{aligned} \quad (78)$$

Allowing for Eq. (60), we obtain

$$\begin{aligned} & \frac{1}{2} \left\{ \partial_r^{\dot{k}} (\beta_5 \Psi^{(1)\dot{n}} + \beta_6 \Psi^{(2)\dot{n}} + \beta_7 \Psi^{(3)\dot{n}} + \beta_8 \Psi^{(4)\dot{n}}) \right. \\ & \left. + \partial_r^{\dot{n}} (\beta_5 \Psi^{(1)\dot{k}} + \beta_6 \Psi^{(2)\dot{k}} + \beta_7 \Psi^{(3)\dot{k}} + \beta_8 \Psi^{(4)\dot{k}}) \right\} + \frac{1}{2} M \left\{ (\sigma^\mu)_{\dot{k}}^{\dot{k}} \Psi_\mu^{\dot{n}} + (\sigma^\mu)_{\dot{n}}^{\dot{n}} \Psi_\mu^{\dot{k}} \right\} = 0. \end{aligned} \quad (79)$$

Let $\dot{A} = \dot{r}$, $D = (kn)$, then from Eq. (77) it follows

$$\begin{aligned} & \frac{1}{2} \left\{ \partial_k^{\dot{r}} (\beta_5 \Psi_n^{(1)} + \beta_6 \Psi_n^{(2)} + \beta_7 \Psi_n^{(3)} + \beta_8 \Psi_n^{(4)}) \right. \\ & \left. + \partial_n^{\dot{r}} (\beta_5 \Psi_k^{(1)} + \beta_6 \Psi_k^{(2)} + \beta_7 \Psi_k^{(3)} + \beta_8 \Psi_k^{(4)}) \right\} + M \Psi_{(kn)}^{\dot{r}} = 0, \end{aligned}$$

$$\begin{aligned} & \frac{1}{2} \left\{ \partial_k^{\dot{r}} (\beta_5 \Psi_n^{(1)} + \beta_6 \Psi_n^{(2)} + \beta_7 \Psi_n^{(3)} + \beta_8 \Psi_n^{(4)}) \right. \\ & \left. + \partial_n^{\dot{r}} (\beta_5 \Psi_k^{(1)} + \beta_6 \Psi_k^{(2)} + \beta_7 \Psi_k^{(3)} + \beta_8 \Psi_k^{(4)}) \right\} + \frac{1}{2} M \left\{ (\sigma^\mu)_{\dot{k}}^{\dot{k}} \Psi_{n\mu} + (\sigma^\mu)_{\dot{n}}^{\dot{n}} \Psi_{k\mu} \right\} = 0. \end{aligned} \quad (80)$$

Let us act on Eq. (79) by $(\sigma^\lambda)_{\dot{k}}^{\dot{k}}$, and on Eq. (80) – by $(\sigma^\lambda)_{\dot{r}}^{\dot{r}}$ after this we convolute the spinor indices; in this way, we arrive at

$$-\frac{1}{2} \left\{ \sigma^{\lambda\dot{k}\dot{r}} \partial_{r\dot{k}} (\beta_5 \Psi^{(1)\dot{n}} + \beta_6 \Psi^{(2)\dot{n}} + \beta_7 \Psi^{(3)\dot{n}} + \beta_8 \Psi^{(4)\dot{n}}) \right.$$

$$\begin{aligned}
& +\partial^{\dot{n}r}(\sigma^\lambda)_{rk}(\beta_5\Psi^{(1)\dot{k}} + \beta_6\Psi^{(2)\dot{k}} + \beta_7\Psi^{(3)\dot{k}} + \beta_8\Psi^{(4)\dot{k}})\Big\} \\
& -\frac{1}{2}M\Big\{(\sigma^\lambda)^{\dot{k}r}(\sigma^\mu)_{rk}\Psi_{\dot{\mu}}^{\dot{n}} + (\sigma^\mu)^{\dot{n}r}(\sigma^\lambda)_{rk}\Psi_{\dot{\mu}}^{\dot{k}}\Big\} = 0, \\
& -\frac{1}{2}\Big\{(\sigma^\lambda)_{kr}\partial^{\dot{r}k}(\beta_5\Psi_n^{(1)} + \beta_6\Psi_n^{(2)} + \beta_7\Psi_n^{(3)} + \beta_8\Psi_n^{(4)}) \\
& +\partial_{\dot{n}r}(\sigma^\lambda)_{rk}(\beta_5\Psi_k^{(1)} + \beta_6\Psi_k^{(2)} + \beta_7\Psi_k^{(3)} + \beta_8\Psi_k^{(4)})\Big\} \\
& -\frac{1}{2}M\Big\{(\sigma^\lambda)_{kr}(\sigma^\mu)^{\dot{r}k}\Psi_{n\mu} + (\sigma^\mu)_{nr}(\sigma^\lambda)^{\dot{r}k}\Psi_{k\mu}\Big\} = 0,
\end{aligned}$$

or

$$\begin{aligned}
& \frac{1}{2}\Big\{\frac{2}{i}\partial_\lambda(\beta_5\Psi^{(1)\dot{n}} + \beta_6\Psi^{(2)\dot{n}} + \beta_7\Psi^{(3)\dot{n}} + \beta_8\Psi^{(4)\dot{n}}) \\
& -\partial_{\dot{n}r}(\sigma^\lambda)_{rk}(\beta_5\Psi^{(1)\dot{k}} + \beta_6\Psi^{(2)\dot{k}} + \beta_7\Psi^{(3)\dot{k}} + \beta_8\Psi^{(4)\dot{k}})\Big\} + \frac{1}{2}M\Big\{2\Psi_{\dot{\lambda}}^{\dot{n}} - (\sigma^\mu)^{\dot{n}r}(\sigma^\lambda)_{rk}\Psi_{\dot{\mu}}^{\dot{k}}\Big\} = 0, \tag{81}
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{2}\Big\{\frac{2}{i}\partial_\lambda(\beta_5\Psi_n^{(1)} + \beta_6\Psi_n^{(2)} + \beta_7\Psi_n^{(3)} + \beta_8\Psi_n^{(4)}) \\
& -\partial_{n\dot{r}}(\sigma^\lambda)^{\dot{r}k}(\beta_5\Psi_k^{(1)} + \beta_6\Psi_k^{(2)} + \beta_7\Psi_k^{(3)} + \beta_8\Psi_k^{(4)})\Big\} + \frac{1}{2}M\Big\{2\Psi_{n\lambda} - (\sigma^\mu)_{nr}(\sigma^\lambda)^{\dot{r}k}\Psi_{k\mu}\Big\} = 0. \tag{82}
\end{aligned}$$

Joining two last equations(81) and (82) into the matrix one, we get

$$\begin{aligned}
& \frac{1}{2}\Big\{\frac{2}{i}\partial_\lambda\Big[\beta_5\begin{vmatrix}\Psi^{(1)\dot{n}} \\ \Psi_n^{(1)}\end{vmatrix} + \beta_6\begin{vmatrix}\Psi^{(2)\dot{n}} \\ \Psi_n^{(2)}\end{vmatrix} + \beta_7\begin{vmatrix}\Psi^{(3)\dot{n}} \\ \Psi_n^{(3)}\end{vmatrix} + \beta_8\begin{vmatrix}\Psi^{(4)\dot{n}} \\ \Psi_n^{(4)}\end{vmatrix}\Big] \\
& \begin{vmatrix}\partial^{\dot{n}r}(\sigma^\lambda)_{rk} & \partial_{n\dot{r}}(\sigma^\lambda)^{\dot{r}k} \\ \Psi_{\dot{\lambda}}^{\dot{n}} & \Psi_{n\lambda}\end{vmatrix} \begin{vmatrix}\Psi^{(1)\dot{k}} \\ \Psi_k^{(1)}\end{vmatrix} + \beta_6\begin{vmatrix}\Psi^{(2)\dot{k}} \\ \Psi_k^{(2)}\end{vmatrix} + \beta_7\begin{vmatrix}\Psi^{(3)\dot{k}} \\ \Psi_k^{(3)}\end{vmatrix} + \beta_8\begin{vmatrix}\Psi^{(4)\dot{k}} \\ \Psi_k^{(4)}\end{vmatrix} \\
& +\frac{1}{2}M\Big\{2\begin{vmatrix}\Psi_{\dot{\lambda}}^{\dot{n}} \\ \Psi_{n\lambda}\end{vmatrix} - \begin{vmatrix}(\sigma^\mu)^{\dot{n}r}(\sigma^\lambda)_{rk} & (\sigma^\mu)_{nr}(\sigma^\lambda)^{\dot{r}k} \\ \Psi_{\dot{\mu}}^{\dot{k}} & \Psi_{k\mu}\end{vmatrix}\Big\} = 0.
\end{aligned}$$

Allowing for (66), we rewrite the last equation in the form

$$\frac{1}{2}\Big\{-i\partial_\lambda + i\gamma_\lambda\hat{\partial} - 2i\partial_\lambda\Big\}(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) + \frac{1}{2}M\Big\{2\Psi_\lambda - \gamma_\lambda(\gamma_\mu\Psi_\mu) + 2\Psi_\lambda\Big\} = 0,$$

or

$$-i(\partial_\lambda - \frac{1}{4}\gamma_\lambda\hat{\partial})(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) + M\Big\{\Psi_\lambda - \frac{1}{4}\gamma_\lambda(\gamma_\mu\Psi_\mu)\Big\} = 0. \tag{83}$$

3. The minimal form of equations

Let us derive the minimal form of the above system (67),(70),(73),(76),(83). From (83) it follows

$$\Psi_\lambda - \frac{1}{4}\gamma_\lambda(\gamma_\mu\Psi_\mu) = \frac{i}{M}(\partial_\lambda - \frac{1}{4}\gamma_\lambda\hat{\partial})(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}). \tag{84}$$

Let us substitute this expression (84) in equations (67), (70), (73), (76). We will take into account the identities

$$\begin{aligned}
\partial_\lambda(\Psi_\lambda - \frac{1}{4}\gamma_\lambda(\gamma_\mu\Psi_\mu)) &= \partial_\lambda\Psi_\lambda - \frac{1}{4}\hat{\partial}(\gamma_\mu\Psi_\mu) = \frac{i}{M}(\square - \frac{1}{4}\hat{\partial}\hat{\partial})(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) \\
&= \frac{i3}{4M}\square(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}),
\end{aligned}$$

so we derive

$$(M + \lambda_1\hat{\partial})\Psi^{(1)} + \frac{3\beta_1}{2M}\square(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0, \tag{85}$$

$$(M + \lambda_2\hat{\partial})\Psi^{(2)} + \frac{3\beta_2}{2M}\square(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0, \tag{86}$$

$$(M + \lambda_3\hat{\partial})\Psi^{(3)} + \frac{3\beta_3}{2M}\square(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0, \tag{87}$$

$$(M + \lambda_4\hat{\partial})\Psi^{(4)} + \frac{3\beta_4}{2M}\square(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0. \tag{88}$$

Let us act on Eq. (85) by operator

$$\beta_5(M + \lambda_2 \hat{\partial})(M + \lambda_3 \hat{\partial})(M + \lambda_4 \hat{\partial});$$

on Eq. (86) – by

$$\beta_6(M + \lambda_1 \hat{\partial})(M + \lambda_3 \hat{\partial})(M + \lambda_4 \hat{\partial});$$

on Eq. (87) by

$$\beta_7(M + \lambda_1 \hat{\partial})(M + \lambda_2 \hat{\partial})(M + \lambda_4 \hat{\partial});$$

on Eq. (88) – by

$$\beta_8(M + \lambda_1 \hat{\partial})(M + \lambda_2 \hat{\partial})(M + \lambda_3 \hat{\partial});$$

and sum the results – this gives

$$\begin{aligned} & \left\{ M^4 + M^3(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)\hat{\partial} + M^2(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)\square \right. \\ & \left. + M(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4)\hat{\partial}\square + \lambda_1\lambda_2\lambda_3\lambda_4\square\square \right\} (\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) \\ & + \frac{3}{2M} \left\{ M^3(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8) \right. \\ & + M^2[(\lambda_2 + \lambda_3 + \lambda_4)\beta_1\beta_5 + (\lambda_1 + \lambda_3 + \lambda_4)\beta_2\beta_6 + (\lambda_1 + \lambda_2 + \lambda_4)\beta_3\beta_7 \\ & + (\lambda_1 + \lambda_2 + \lambda_3)\beta_4\beta_8]\hat{\partial} + M[\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) \\ & + \beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) + \beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3)]\square \\ & \left. + [\lambda_1\lambda_2(\lambda_4\beta_3\beta_7 + \lambda_3\beta_4\beta_8) + \lambda_3\lambda_4(\lambda_2\beta_1\beta_5 + \lambda_1\beta_2\beta_6)]\hat{\partial}\square \right\} \\ & \times (\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0, \end{aligned}$$

or

$$\begin{aligned} & \left\{ M^4 + M^3(\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4)\hat{\partial} + M^2(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) \right. \\ & \left. + \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)\square \right. \\ & + M \left[(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) + \frac{3}{2}((\lambda_2 + \lambda_3 + \lambda_4)\beta_1\beta_5 \right. \\ & + (\lambda_1 + \lambda_3 + \lambda_4)\beta_2\beta_6 + (\lambda_1 + \lambda_2 + \lambda_4)\beta_3\beta_7 + (\lambda_1 + \lambda_2 + \lambda_3)\beta_4\beta_8) \Big] \hat{\partial}\square \\ & + \left[\lambda_1\lambda_2\lambda_3\lambda_4 + \frac{3}{2}(\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) \right. \\ & + \beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) + \beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3)) \Big] \square\square \\ & \left. + [\lambda_1\lambda_2(\lambda_4\beta_3\beta_7 + \lambda_3\beta_4\beta_8) + \lambda_3\lambda_4(\lambda_2\beta_1\beta_5 + \lambda_1\beta_2\beta_6)]\hat{\partial}\square\square \right\} \\ & \times (\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0. \end{aligned} \tag{89}$$

Taking in mind restrictions (16), (20), (28), (37), (41) for parameters λ_i , β_k , we derive

$$(M + \hat{\partial}) \left(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)} \right) = 0, \tag{90}$$

which is the minimal Dirac equation for a free p[article with spin 1/2, with respect to the following bispinor

$$\Phi = \beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}. \tag{91}$$

4. Equation in presence of electromagnetic fields

We start with the equations in presence of external electromagnetic fields ($D_\mu = \partial_\mu - ieA_\mu$)

$$(M + \lambda_1 \hat{D})\Psi^{(1)} - i2\beta_1 \left\{ D_\mu \Psi_\mu - \frac{1}{4} \hat{D}(\gamma_\mu \Psi_\mu) \right\} = 0, \quad (92)$$

$$(M + \lambda_2 \hat{D})\Psi^{(2)} - i2\beta_2 \left\{ D_\mu \Psi_\mu - \frac{1}{4} \hat{D}(\gamma_\mu \Psi_\mu) \right\} = 0, \quad (93)$$

$$(M + \lambda_3 \hat{D})\Psi^{(3)} - i2\beta_3 \left\{ D_\mu \Psi_\mu - \frac{1}{4} \hat{D}(\gamma_\mu \Psi_\mu) \right\} = 0, \quad (94)$$

$$(M + \lambda_4 \hat{D})\Psi^{(4)} - i2\beta_4 \left\{ D_\mu \Psi_\mu - \frac{1}{4} \hat{D}(\gamma_\mu \Psi_\mu) \right\} = 0, \quad (95)$$

$$-i(D_\lambda - \frac{1}{4}\gamma_\lambda \hat{D})(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) + M \left\{ \Psi_\lambda - \frac{1}{4}\gamma_\lambda(\gamma_\mu \Psi_\mu) \right\} = 0. \quad (96)$$

From Eq. (96), it follows

$$\Psi_\lambda - \frac{1}{4}\gamma_\lambda(\gamma_\mu \Psi_\mu) = \frac{i}{M}(D_\lambda - \frac{1}{4}\gamma_\lambda \hat{D})(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}).$$

Correspondingly, we obtain

$$D_\mu \Psi_\mu - \frac{1}{4} \hat{D}(\gamma_\mu \Psi_\mu) = \frac{i}{M}(D^2 - \frac{1}{4} \hat{D} \hat{D})(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}). \quad (97)$$

Taking into account this identity (97) in equations (92)–(95), we derive

$$(M + \lambda_1 D)\Psi^{(1)} + \frac{2\beta_1}{M}(D^2 - \frac{1}{4} \hat{D} \hat{D})(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0, \quad (98)$$

$$(M + \lambda_2 D)\Psi^{(2)} + \frac{2\beta_2}{M}(D^2 - \frac{1}{4} \hat{D} \hat{D})(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0, \quad (99)$$

$$(M + \lambda_3 D)\Psi^{(3)} + \frac{2\beta_3}{M}(D^2 - \frac{1}{4} \hat{D} \hat{D})(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0, \quad (100)$$

$$(M + \lambda_4 D)\Psi^{(4)} + \frac{2\beta_4}{M}(D^2 - \frac{1}{4} \hat{D} \hat{D})(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0. \quad (101)$$

Let us act on Eq. (98) by operator

$$\beta_5(M + \lambda_2 \hat{D})(M + \lambda_3 \hat{D})(M + \lambda_4 \hat{D}),$$

on Eq. (99) – by

$$\beta_6(M + \lambda_1 \hat{D})(M + \lambda_3 \hat{D})(M + \lambda_4 \hat{D}),$$

in Eq. (100) – by

$$\beta_7(M + \lambda_1 \hat{D})(M + \lambda_2 \hat{D})(M + \lambda_4 \hat{D}),$$

on Eq. (101) – by

$$\beta_8(M + \lambda_1 \hat{D})(M + \lambda_2 \hat{D})(M + \lambda_3 \hat{D}),$$

and sum the results; so we derive

$$\begin{aligned} & \left\{ M^4 + M^3 \hat{D} + M^2(\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4) \hat{D} \hat{D} \right. \\ & \quad + M(\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4) \hat{D} \hat{D} \hat{D} \\ & \quad \left. + \lambda_1 \lambda_2 \lambda_3 \lambda_4 \hat{D} \hat{D} \hat{D} \hat{D} \right\} (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) \\ & + \frac{2}{M} \left\{ \beta_1 \beta_5 [M^3 + M^2(\lambda_2 + \lambda_3 + \lambda_4) \hat{D} + M(\lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4) \hat{D} \hat{D} + \lambda_2 \lambda_3 \lambda_4 \hat{D} \hat{D} \hat{D}] \right. \\ & \quad + \beta_2 \beta_6 [M^3 + M^2(\lambda_1 + \lambda_3 + \lambda_4) \hat{D} + M(\lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_3 \lambda_4) \hat{D} \hat{D} + \lambda_1 \lambda_3 \lambda_4 \hat{D} \hat{D} \hat{D}] \\ & \quad \left. + \beta_3 \beta_7 [M^3 + M^2(\lambda_1 + \lambda_2 + \lambda_4) \hat{D} + M(\lambda_1 \lambda_2 + \lambda_1 \lambda_4 + \lambda_2 \lambda_4) \hat{D} \hat{D} + \lambda_1 \lambda_2 \lambda_4 \hat{D} \hat{D} \hat{D}] \right\} \end{aligned}$$

$$+ \beta_4 \beta_8 [M^3 + M^2(\lambda_1 + \lambda_2 + \lambda_3)\hat{D} + M(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3)\hat{D}\hat{D} + \lambda_1\lambda_2\lambda_3\hat{D}\hat{D}\hat{D}] \times (D^2 - \frac{1}{4}\hat{D}\hat{D}) (\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0. \quad (102)$$

We will use the identity $\hat{D}\hat{D} = D^2 - ieF_{[\mu\nu]}I_{[\mu\nu]}$, so that

$$D^2 = \hat{D}\hat{D} + ieF_{[\mu\nu]}I_{[\mu\nu]}; \quad (103)$$

where $F_{[\mu\nu]} = \partial_\mu A_\nu - \partial_\nu A_\mu$ is the electromagnetic tensor, and $I_{[\mu\nu]} = \frac{1}{4}(\gamma_\mu\gamma_\nu - \gamma_\nu\gamma_\mu)$. Allowing for relations (103), we transform Eq. (102) to the form

$$\begin{aligned} & \left\{ M^4 + M^3\hat{D} + M^2[(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) \right. \\ & \quad + \frac{3}{2}(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)]\hat{D}\hat{D} \\ & \quad + M[(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4) \\ & \quad + \frac{3}{2}\beta_1\beta_5(\lambda_2 + \lambda_3 + \lambda_4) + \frac{3}{2}\beta_2\beta_6(\lambda_1 + \lambda_3 + \lambda_4) + \frac{3}{2}\beta_3\beta_7(\lambda_1 + \lambda_2 + \lambda_4) + \frac{3}{2}\beta_4\beta_8(\lambda_1 + \lambda_2 + \lambda_3)]\hat{D}\hat{D}\hat{D} \\ & \quad + \left[\lambda_1\lambda_2\lambda_3\lambda_4 + \frac{3}{2}\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + \frac{3}{2}\beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) \right. \\ & \quad \left. + \frac{3}{2}\beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \frac{3}{2}\beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3) \right] \hat{D}\hat{D}\hat{D}\hat{D} \\ & \left. + \frac{3}{2M}[\beta_1\beta_5\lambda_2\lambda_3\lambda_4 + \beta_2\beta_6\lambda_1\lambda_3\lambda_4 + \beta_3\beta_7\lambda_1\lambda_2\lambda_4 + \beta_4\beta_8\lambda_1\lambda_2\lambda_3]\hat{D}\hat{D}\hat{D}\hat{D}\hat{D} \right\} (\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) \\ & \quad + \frac{2ie}{M} \left\{ M^3(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)F_{[\mu\nu]}I_{[\mu\nu]} \right. \\ & \quad + M^2[\beta_1\beta_5(\lambda_2 + \lambda_3 + \lambda_4) + \frac{3}{2}\beta_2\beta_6(\lambda_1 + \lambda_3 + \lambda_4) + \frac{3}{2}\beta_3\beta_7(\lambda_1 + \lambda_2 + \lambda_4) + \frac{3}{2}\beta_4\beta_8(\lambda_1 + \lambda_2 + \lambda_3)]\hat{D}F_{[\mu\nu]}I_{[\mu\nu]} \\ & \quad + M[\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + \beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) \\ & \quad + \beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3)]\hat{D}\hat{D}F_{[\mu\nu]}I_{[\mu\nu]} \\ & \quad \left. + [\beta_1\beta_5\lambda_2\lambda_3\lambda_4 + \beta_2\beta_6\lambda_1\lambda_3\lambda_4 + \beta_3\beta_7\lambda_1\lambda_2\lambda_4 + \beta_4\beta_8\lambda_1\lambda_2\lambda_3]\hat{D}\hat{D}\hat{D}F_{[\mu\nu]}I_{[\mu\nu]} \right\} \\ & \quad \times (\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)}) = 0. \end{aligned}$$

Taking in mind restrictions (20), (28), (41), (37), and (91), we derive

$$\begin{aligned} & \left\{ M^4 + M^3\hat{D} + 2ieM^2(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8)F_{[\mu\nu]}I_{[\mu\nu]} \right. \\ & \quad + 2ieM[\beta_1\beta_5(\lambda_2 + \lambda_3 + \lambda_4) + \beta_2\beta_6(\lambda_1 + \lambda_3 + \lambda_4) + \beta_3\beta_7(\lambda_1 + \lambda_2 + \lambda_4) + \beta_4\beta_8(\lambda_1 + \lambda_2 + \lambda_3)]\hat{D}F_{[\mu\nu]}I_{[\mu\nu]} \\ & \quad + 2ie[\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + \beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) \\ & \quad + \beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3)]\hat{D}\hat{D}F_{[\mu\nu]}I_{[\mu\nu]} \left. \right\} \Phi = 0, \quad (104) \end{aligned}$$

or differently

$$\begin{aligned} & M + \hat{D} - ie\frac{4}{3}\left[\frac{1}{M}(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)F_{[\mu\nu]}I_{[\mu\nu]} \right. \\ & \quad \left. + \frac{1}{M^2}(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4)\hat{D}F_{[\mu\nu]}I_{[\mu\nu]} + \frac{1}{M^3}\lambda_1\lambda_2\lambda_3\lambda_4\hat{D}\hat{D}F_{[\mu\nu]}I_{[\mu\nu]} \right] \Phi = 0. \quad (105) \end{aligned}$$

Allowing for the identities

$$\hat{D}\hat{D} = D^2 - ieF_{[\rho\lambda]}I_{[\rho\lambda]},$$

$$\gamma_\mu\gamma_\nu\gamma_\rho = \delta_{\mu\nu}\gamma_\rho - \delta_{\mu\rho}\gamma_\nu + \delta_{\nu\rho}\gamma_\mu + \epsilon_{\mu\nu\rho\eta}\gamma_5\gamma_\eta, \quad \gamma_5\gamma_\mu\gamma_\nu = \delta_{\mu\nu}\gamma_5 - \frac{1}{2}\epsilon_{\mu\nu\rho\lambda}\gamma_\rho\gamma_\lambda.$$

we get

$$\begin{aligned} & \frac{1}{M^3}\lambda_1\lambda_2\lambda_3\lambda_4 \left\{ D^2F_{[\mu\nu]}I_{[\mu\nu]}\Phi - ieF_{[\rho\lambda]}F_{[\mu\nu]}I_{[\rho\lambda]}I_{[\mu\nu]}\Phi \right\} = \frac{1}{M^3}\lambda_1\lambda_2\lambda_3\lambda_4 D^2F_{[\mu\nu]}I_{[\mu\nu]}\Phi, \\ & - \frac{ie}{M^3}\lambda_1\lambda_2\lambda_3\lambda_4 F_{[\rho\lambda]}F_{[\mu\nu]}I_{[\rho\lambda]}I_{[\mu\nu]} = \frac{1}{M^3}\lambda_1\lambda_2\lambda_3\lambda_4 D^2F_{[\mu\nu]}I_{[\mu\nu]}\Phi, \\ & - \frac{ie}{M^3}\lambda_1\lambda_2\lambda_3\lambda_4 \left\{ F_{[\mu\nu]}I_{[\nu\rho]}\gamma_\mu\gamma_\rho + \frac{1}{4}\epsilon_{[\rho\lambda]}F_{[\mu\nu]}\gamma_5 - \frac{1}{2}F_{[\mu\nu]}F_{[\nu\mu]} \right\} \Phi. \end{aligned}$$

Therefore, Eq. (124) may be presented as

$$\begin{aligned} & \left\{ M + \gamma_\sigma D_\sigma - \frac{ie}{M} \frac{4}{3} \left(\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4 \right) F_{[\mu\nu]} I_{[\mu\nu]} \right. \\ & \quad \left. - \frac{ie}{M^2} \frac{4}{3} \left(\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4 \right) \hat{D} F_{[\mu\nu]} I_{\mu\nu} \right. \\ & \quad \left. - \frac{ie}{M^3} \frac{4}{3} \lambda_1 \lambda_2 \lambda_3 \lambda_4 \left(D^2 F_{[\mu\nu]} I_{[\mu\nu]} + F_{[\mu\nu]} F_{[\nu\rho]} \gamma_\mu \gamma_\rho + \frac{1}{2} F_{[\mu\nu]} F_{[\mu\nu]} + \frac{1}{4} \epsilon_{\rho\lambda\mu\nu} F_{[\rho\lambda]} F_{[\mu\nu]} \gamma_5 \right) \right\} \Phi = 0. \end{aligned} \quad (106)$$

Let us introduce the notations

$$\begin{aligned} \mu &= \frac{4}{3} \left(\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4 \right), \\ \sigma &= \frac{4}{3} \left(\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4 \right), \quad \eta = \frac{4}{3} \lambda_1 \lambda_2 \lambda_3 \lambda_4; \end{aligned}$$

then the above equation reads

$$\begin{aligned} & \left\{ M + \gamma_\rho D_\rho - \frac{ie}{M} \mu F_{\mu\beta} I_{\mu\beta} - \frac{ie}{M^2} \sigma \hat{D} F_{\mu\beta} I_{\mu\beta} \right. \\ & \quad \left. - \frac{ie}{M^3} \eta \left(D^2 F_{\alpha\beta} I_{\alpha\beta} + F_{\alpha\beta} F_{\beta\rho} \gamma_\alpha \gamma_\rho + \frac{1}{2} F_{\alpha\beta} F_{[\alpha\beta]} + \frac{1}{4} \epsilon_{\rho\lambda\alpha\beta} F_{\rho\lambda} F_{\alpha\beta} \gamma_5 \right) \right\} \Phi = 0. \end{aligned} \quad (107)$$

This equation should be considered as related to a spin 1/2 particle which in addition to electric charge has three other electromagnetic characteristics μ, σ, γ . The μ corresponds to the anomalous magnetic moment; the σ corresponds to the polarizability; and the parameter γ should be associated with a certain intrinsic structure which assumes rather complicated additional interaction with the external electromagnetic field.

5. Equation in presence of gravitational fields

Assuming the use of relativistic interval in explicit real-valued form (Minkowski space with the signature $(+, -, -, -)$)

$$dS^2 = c^2 dt^2 - dx^2 - dy^2 - dz^2, \quad (t, x, y, z) = (x^a), \quad a = 0, 1, 2, 3$$

one should use the basic equation in the form

$$\begin{aligned} & \left\{ i\gamma^\rho D_\rho - M + \frac{e\mu}{M} F_{\alpha\beta} J^{\alpha\beta} + \frac{e\sigma}{M^2} \hat{D} F_{\alpha\beta} J^{\alpha\beta} \right. \\ & \quad \left. + \frac{e\eta}{M^3} \left(D^2 F_{\alpha\beta} J^{\alpha\beta} + \gamma^\alpha F_{\alpha\beta} F_{\beta\rho} \gamma^\rho + \frac{1}{2} F_{\alpha\beta} F^{\alpha\beta} + \frac{1}{4} \gamma_5 \epsilon^{\rho\lambda\alpha\beta} F_{\rho\lambda} F_{\alpha\beta} \right) \right\} \Psi = 0; \end{aligned} \quad (108)$$

recall the notations

$$D_\rho = \partial_\rho + ieA_\rho, \quad \hat{D} = \gamma^\rho D_\rho, \quad D^2 = g^{\alpha\beta} D_\alpha D_\beta, \quad J^{\alpha\beta} = \frac{1}{4} (\gamma^\alpha \gamma^\beta - \gamma^\beta \gamma^\alpha).$$

Let us extend this approach to curved space-time background. We should start with the modified system of the first order equations

$$(M + \lambda_1 \hat{D}) \Psi^{(1)} - 2i\beta_1 \left(D^\mu \Psi_\mu - \frac{1}{4} \hat{D} (\gamma^\mu \Psi_\mu) \right) = 0, \quad (109)$$

$$(M + \lambda_2 \hat{D}) \Psi^{(2)} - 2i\beta_2 \left(D^\mu \Psi_\mu - \frac{1}{4} \hat{D} (\gamma^\mu \Psi_\mu) \right) = 0, \quad (110)$$

$$(M + \lambda_3 \hat{D}) \Psi^{(3)} - 2i\beta_3 \left(D^\mu \Psi_\mu - \frac{1}{4} \hat{D} (\gamma^\mu \Psi_\mu) \right) = 0, \quad (111)$$

$$(M + \lambda_4 \hat{D}) \Psi^{(4)} - 2i\beta_4 \left(D^\mu \Psi_\mu - \frac{1}{4} \hat{D} (\gamma^\mu \Psi_\mu) \right) = 0, \quad (112)$$

$$-i(D_\lambda - \frac{1}{4} \gamma_\lambda \hat{D}) (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) + M \left\{ \Psi_\lambda - \frac{1}{4} \gamma_\lambda (\gamma^\mu \Psi_\mu) \right\} = 0. \quad (113)$$

In this system, $\Psi^{(a)} (a = 1, 2, 3, 4)$ are covariant bispinors, and Ψ_μ is a covariant vector-bispinor. We apply the generalized derivative D_μ , it takes into account the presence of electromagnetic fields and any

curved space-time background; the symbol ∇_μ designates the covariant derivative; the symbol $\Gamma_\mu(x)$ stands for the bispinor connection; $\Gamma_\mu(x)$ represents the local Dirac matrices:

$$\begin{aligned} D_\mu &= \nabla_\mu - ieA_\mu(x) + \Gamma_\mu(x), \quad \hat{D} = \gamma^\mu D_\mu = \gamma^\mu(x) \left(\nabla_\mu - ieA_\mu + \Gamma_\mu \right), \\ \hat{D}\hat{D} &= (\gamma^\alpha D_\alpha)(\gamma^\beta D_\beta) = D_\alpha \frac{\gamma^\alpha \gamma^\beta + \gamma^\beta \gamma^\alpha}{2} D_\beta + D_\alpha \frac{\gamma^\alpha \gamma^\beta - \gamma^\beta \gamma^\alpha}{2} D_\beta \\ &= g^{\alpha\beta}(x) D_\alpha D_\beta = D^\alpha D_\alpha + J^{\alpha\beta}(x) [D_\alpha, D_\beta]_- = D^2 + \sigma^{\alpha\beta}(x) M_{\alpha\beta}(x), \\ D^2 &= D^\alpha D_\alpha, \quad M_{\alpha\beta}(x) = [D_\alpha, D_\beta]_-, \quad J^{\alpha\beta}(x) = \frac{\gamma^\alpha(x) \gamma^\beta(x) - \gamma^\beta(x) \gamma^\alpha(x)}{4}. \end{aligned} \quad (114)$$

From Eq. (113), it follows

$$\Psi_\lambda - \frac{1}{4} \gamma_\lambda (\gamma^\mu \Psi_\mu) = \frac{i}{M} \left(D_\lambda - \frac{1}{4} \gamma_\lambda \hat{D} \right) \left(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)} \right).$$

Correspondingly, acting on this relation by D_λ , we obtain

$$D^\lambda \Psi_\lambda - \frac{1}{4} \hat{D} (\gamma^\mu \Psi_\mu) = \frac{i}{M} \left(D^2 - \frac{1}{4} \hat{D} \hat{D} \right) \left(\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)} \right). \quad (115)$$

Taking into account the last identity in equations (109)–(112), we derive

$$(M + \lambda_1 \hat{D}) \Psi^{(1)} + \frac{2\beta_1}{M} \left(D^2 - \frac{1}{4} \hat{D} \hat{D} \right) (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0, \quad (116)$$

$$(M + \lambda_2 \hat{D}) \Psi^{(2)} + \frac{2\beta_2}{M} \left(D^2 - \frac{1}{4} \hat{D} \hat{D} \right) (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0, \quad (117)$$

$$(M + \lambda_3 \hat{D}) \Psi^{(3)} + \frac{2\beta_3}{M} \left(D^2 - \frac{1}{4} \hat{D} \hat{D} \right) (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0, \quad (118)$$

$$(M + \lambda_4 \hat{D}) \Psi^{(4)} + \frac{2\beta_4}{M} \left(D^2 - \frac{1}{4} \hat{D} \hat{D} \right) (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0. \quad (119)$$

Let us act:

on Eq. (116) by operator

$$\beta_5 (M + \lambda_2 \hat{D}) (M + \lambda_3 \hat{D}) (M + \lambda_4 \hat{D}),$$

on Eq. (117) – by

$$\beta_6 (M + \lambda_1 \hat{D}) (M + \lambda_3 \hat{D}) (M + \lambda_4 \hat{D}),$$

in Eq. (118) – by

$$\beta_7 (M + \lambda_1 \hat{D}) (M + \lambda_2 \hat{D}) (M + \lambda_4 \hat{D}),$$

on Eq. (119) – by

$$\beta_8 (M + \lambda_1 \hat{D}) (M + \lambda_2 \hat{D}) (M + \lambda_3 \hat{D}),$$

and sum the results; so we derive

$$\begin{aligned} & \left\{ M^4 + M^3 \hat{D} + M^2 (\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4) \hat{D} \hat{D} \right. \\ & + M (\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4) \hat{D} \hat{D} \hat{D} + \lambda_1 \lambda_2 \lambda_3 \lambda_4 \hat{D} \hat{D} \hat{D} \hat{D} \left. \right\} (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) \\ & + \frac{2}{M} \left\{ \beta_1 \beta_5 [M^3 + M^2 (\lambda_2 + \lambda_3 + \lambda_4) \hat{D} + M (\lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4) \hat{D} \hat{D} + \lambda_2 \lambda_3 \lambda_4 \hat{D} \hat{D} \hat{D}] \right. \\ & + \beta_2 \beta_6 [M^3 + M^2 (\lambda_1 + \lambda_3 + \lambda_4) \hat{D} + M (\lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_3 \lambda_4) \hat{D} \hat{D} + \lambda_1 \lambda_3 \lambda_4 \hat{D} \hat{D} \hat{D}] \\ & + \beta_3 \beta_7 [M^3 + M^2 (\lambda_1 + \lambda_2 + \lambda_4) \hat{D} + M (\lambda_1 \lambda_2 + \lambda_1 \lambda_4 + \lambda_2 \lambda_4) \hat{D} \hat{D} + \lambda_1 \lambda_2 \lambda_4 \hat{D} \hat{D} \hat{D}] \\ & + \beta_4 \beta_8 [M^3 + M^2 (\lambda_1 + \lambda_2 + \lambda_3) \hat{D} + M (\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_2 \lambda_3) \hat{D} \hat{D} + \lambda_1 \lambda_2 \lambda_3 \hat{D} \hat{D} \hat{D}] \left. \right\} \\ & \times \left(D^2 - \frac{1}{4} \hat{D} \hat{D} \right) (\beta_5 \Psi^{(1)} + \beta_6 \Psi^{(2)} + \beta_7 \Psi^{(3)} + \beta_8 \Psi^{(4)}) = 0. \end{aligned} \quad (120)$$

We will use the identity (see (114))

$$\hat{D}\hat{D} = D^2 + J^{\alpha\beta}(x) M_{\alpha\beta}(x) \implies D^2 = \hat{D}\hat{D} - M_{\alpha\beta} J^{\alpha\beta}; \quad (121)$$

instead of relation $D^2 = \hat{D}\hat{D} + ieF_{\alpha\beta}J^{\alpha\beta}$ in presence of only electromagnetic field. Therefore, we do not need to repeat the calculation in Minkovski space, it suffices the formal change

$$ieF_{\alpha\beta} \rightarrow -M_{\alpha\beta}.$$

Allowing for relations (120), we transform Eq. (120) to the form

$$\begin{aligned} & \left\{ M^4 + M^3\hat{D} + M^2 \left(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4 \right) \hat{D}^2 \right. \\ & + M \left(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4 \right) \hat{D}^3 + \lambda_1\lambda_2\lambda_3\lambda_4 \hat{D}^4 \left. \right\} \left(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)} \right) \\ & + \frac{2}{M} \left\{ \beta_1\beta_5 \left[M^3 + M^2(\lambda_2 + \lambda_3 + \lambda_4)\hat{D} + M(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4)\hat{D}^2 + \lambda_2\lambda_3\lambda_4\hat{D}^3 \right] \right. \\ & + \beta_2\beta_6 \left[M^3 + M^2(\lambda_1 + \lambda_3 + \lambda_4)\hat{D} + M(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4)\hat{D}^2 + \lambda_1\lambda_3\lambda_4\hat{D}^3 \right] \\ & + \beta_3\beta_7 \left[M^3 + M^2(\lambda_1 + \lambda_2 + \lambda_4)\hat{D} + M(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4)\hat{D}^2 + \lambda_1\lambda_2\lambda_4\hat{D}^3 \right] \\ & + \beta_4\beta_8 \left[M^3 + M^2(\lambda_1 + \lambda_2 + \lambda_3)\hat{D} + M(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3)\hat{D}^2 + \lambda_1\lambda_2\lambda_3\hat{D}^3 \right] \left. \right\} \\ & \times \left(\frac{3}{4}\hat{D}^2 - \sigma_{\alpha\beta}J^{\alpha\beta} \right) \left(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)} \right) = 0. \end{aligned} \quad (122)$$

Now, we are grouping the terms with respect to

$$\hat{D}, \quad \hat{D}^2, \quad \hat{D}^3, \quad \hat{D}^4, \quad \hat{D}^5;$$

in this way we obtain

$$\begin{aligned} & \left\{ M^4 + M^3\hat{D} + M^2 \left[\left(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4 \right) \frac{3}{2} \left(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8 \right) \right] \hat{D}^2 \right. \\ & + M \left[\left(\lambda_1\lambda_2\lambda_3 + \lambda_1\lambda_2\lambda_4 + \lambda_1\lambda_3\lambda_4 + \lambda_2\lambda_3\lambda_4 \right) + \frac{3}{2}\beta_1\beta_5(\lambda_2 + \lambda_3 + \lambda_4) \right. \\ & + \frac{3}{2}\beta_2\beta_6(\lambda_1 + \lambda_3 + \lambda_4) + \frac{3}{2}\beta_3\beta_7(\lambda_1 + \lambda_2 + \lambda_4) + \frac{3}{2}\beta_4\beta_8(\lambda_1 + \lambda_2 + \lambda_3) \left. \right] \hat{D}^3 \\ & + \left[\lambda_1\lambda_2\lambda_3\lambda_4 + \frac{3}{2}\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + \frac{3}{2}\beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) \right. \\ & + \frac{3}{2}\beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \frac{3}{2}\beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3) \left. \right] \hat{D}^4 \\ & + \frac{3}{2M} \left[\beta_1\beta_5\lambda_2\lambda_3\lambda_4 + \beta_2\beta_6\lambda_1\lambda_3\lambda_4 + \beta_3\beta_7\lambda_1\lambda_2\lambda_4 + \beta_4\beta_8\lambda_1\lambda_2\lambda_3 \right] \hat{D}^5 \left. \right\} \left(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)} \right) \\ & - \frac{2}{M} \left\{ M^3 \left(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8 \right) M_{\alpha\beta}J^{\alpha\beta} \right. \\ & + M^2 \left[\beta_1\beta_5(\lambda_2 + \lambda_3 + \lambda_4) + \frac{3}{2}\beta_2\beta_6(\lambda_1 + \lambda_3 + \lambda_4) + \frac{3}{2}\beta_3\beta_7(\lambda_1 + \lambda_2 + \lambda_4) + \frac{3}{2}\beta_4\beta_8(\lambda_1 + \lambda_2 + \lambda_3) \right] \hat{D} M_{\alpha\beta}J^{\alpha\beta} \\ & + M \left[\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + \beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) \right. \\ & + \beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3) \left. \right] \hat{D}^2 M_{\alpha\beta}J^{\alpha\beta} \\ & + \left[\beta_1\beta_5\lambda_2\lambda_3\lambda_4 + \beta_2\beta_6\lambda_1\lambda_3\lambda_4 + \beta_3\beta_7\lambda_1\lambda_2\lambda_4 + \beta_4\beta_8\lambda_1\lambda_2\lambda_3 \right] \hat{D}^3 M_{\alpha\beta}J^{\alpha\beta} \left. \right\} \\ & \times \left(\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)} \right) = 0. \end{aligned}$$

Taking into account restrictions on parameters (20), (28), (41), (37), and also the notation (91)

$$\beta_5\Psi^{(1)} + \beta_6\Psi^{(2)} + \beta_7\Psi^{(3)} + \beta_8\Psi^{(4)} = \Phi,$$

we derive

$$\begin{aligned} & \left\{ M^4 + M^3\hat{D} - 2M^2 \left(\beta_1\beta_5 + \beta_2\beta_6 + \beta_3\beta_7 + \beta_4\beta_8 \right) M_{\alpha\beta}J^{\alpha\beta} \right. \\ & - 2M \left[\beta_1\beta_5(\lambda_2 + \lambda_3 + \lambda_4) + \beta_2\beta_6(\lambda_1 + \lambda_3 + \lambda_4) + \beta_3\beta_7(\lambda_1 + \lambda_2 + \lambda_4) + \beta_4\beta_8(\lambda_1 + \lambda_2 + \lambda_3) \right] \hat{D} M_{\alpha\beta}J^{\alpha\beta} \\ & - 2 \left[\beta_1\beta_5(\lambda_2\lambda_3 + \lambda_2\lambda_4 + \lambda_3\lambda_4) + \beta_2\beta_6(\lambda_1\lambda_3 + \lambda_1\lambda_4 + \lambda_3\lambda_4) \right. \\ & + \beta_3\beta_7(\lambda_1\lambda_2 + \lambda_1\lambda_4 + \lambda_2\lambda_4) + \beta_4\beta_8(\lambda_1\lambda_2 + \lambda_1\lambda_3 + \lambda_2\lambda_3) \left. \right] \hat{D} \hat{D} M_{\alpha\beta}J^{\alpha\beta} \left. \right\} \Phi = 0, \end{aligned} \quad (123)$$

or differently

$$\left\{ M + \hat{D} + \frac{4}{3} \left[\frac{1}{M} \left(\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4 \right) M_{\alpha\beta} J^{\alpha\beta} \right. \right. \\ \left. \left. + \frac{1}{M^2} \left(\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4 \right) \hat{D} M_{\alpha\beta} J^{\alpha\beta} + \frac{1}{M^3} \lambda_1 \lambda_2 \lambda_3 \lambda_4 \hat{D}^2 M_{\alpha\beta} J^{\alpha\beta} \right] \right\} \Phi = 0.$$

Allowing for the identities $\hat{D}^2 = (D^2 + J^{\rho\sigma} M_{\rho\sigma})$, the last equation transforms to other form

$$\left\{ M + \hat{D} + \frac{1}{M} \frac{4}{3} \left(\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4 \right) J^{\alpha\beta} M_{\alpha\beta} \right. \\ \left. + \frac{1}{M^2} \frac{4}{3} \left(\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4 \right) \hat{D} J^{\alpha\beta} M_{\alpha\beta} + \frac{1}{M^3} \frac{4}{3} \lambda_1 \lambda_2 \lambda_3 \lambda_4 (D^2 + J^{\rho\sigma} M_{\rho\sigma}) J^{\alpha\beta} M_{\alpha\beta} \right\} \Phi = 0,$$

or differently

$$\left\{ M + \hat{D} + \frac{1}{M} \frac{4}{3} \left(\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_1 \lambda_4 + \lambda_2 \lambda_3 + \lambda_2 \lambda_4 + \lambda_3 \lambda_4 \right) J^{\alpha\beta} M_{\alpha\beta} \right. \\ + \frac{1}{M^2} \frac{4}{3} \left(\lambda_1 \lambda_2 \lambda_3 + \lambda_1 \lambda_2 \lambda_4 + \lambda_1 \lambda_3 \lambda_4 + \lambda_2 \lambda_3 \lambda_4 \right) \hat{D} J^{\alpha\beta} M_{\alpha\beta} \\ \left. + \frac{1}{M^3} \frac{4}{3} \lambda_1 \lambda_2 \lambda_3 \lambda_4 \left(D^2 J^{\alpha\beta} M_{\alpha\beta} + J^{\rho\sigma} J^{\alpha\beta} M_{\rho\sigma} M_{\alpha\beta} \right) \right\} \Phi = 0. \quad (124)$$

Further, we will work with the short presentation of the above equation

$$\left\{ M + \hat{D} + \frac{1}{M} \mu j^{\alpha\beta} M_{\alpha\beta} + \frac{1}{M^2} \sigma \hat{D} j^{\alpha\beta} M_{\alpha\beta} \right. \\ \left. + \frac{1}{M^3} \eta \left(D^2 j^{\alpha\beta} M_{\alpha\beta} + j^{\rho\sigma} j^{\alpha\beta} M_{\rho\sigma} M_{\alpha\beta} \right) \right\} \Phi = 0. \quad (125)$$

Let us detail the commutator $M_{\alpha\beta}$:

$$D_\alpha = \nabla_\alpha + ieA_\mu(x) + \Gamma_\alpha(x), \quad M_{\alpha\beta} = D_\alpha D_\beta - D_\beta D_\alpha;$$

it is (see in [3])

$$M_{\alpha\beta} \Psi = (D_\alpha D_\beta - D_\beta D_\alpha) \Psi = (ieF_{\alpha\beta} + \frac{1}{2} j^{\nu\rho} R_{\nu\rho\alpha\beta}) \Psi. \quad (126)$$

Let us consider the term (see in [3])

$$\frac{\mu}{M} j^{\alpha\beta} M_{\alpha\beta} \Psi = \frac{\mu}{M} j^{\alpha\beta} \left(ieF_{\alpha\beta} + \frac{1}{2} j^{\nu\rho} R_{\nu\rho\alpha\beta} \right) \Psi \\ = \frac{\mu}{M} \left(ieF_{\alpha\beta} j^{\alpha\beta} + \frac{1}{2} j^{\alpha\beta} j^{\nu\rho} R_{\nu\rho\alpha\beta} \right) \Psi = \frac{\mu}{M} \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R(x) \right) \Psi. \quad (127)$$

Consider the term

$$\frac{\sigma}{M^2} \hat{D} j^{\alpha\beta} M_{\alpha\beta} \Psi = \frac{1}{M^2} \sigma (\gamma^\rho D_\rho) \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R(x) \right) \Psi. \quad (128)$$

Consider the term

$$\frac{\eta}{M^3} D^2 j^{\alpha\beta} M_{\alpha\beta} \Psi = \frac{1}{M^3} \eta g^{\rho\sigma}(x) D_\rho D_\sigma \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R(x) \right) \Psi. \quad (129)$$

Now we are to consider the term

$$A = \frac{\eta}{M^3} j^{\rho\sigma} j^{\alpha\beta} M_{\rho\sigma} M_{\alpha\beta} = \frac{\eta}{M^3} j^{\rho\sigma} \left\{ (j^{\alpha\beta} M_{\rho\sigma} - M_{\rho\sigma} j^{\alpha\beta}) + M_{\rho\sigma} j^{\alpha\beta} \right\} M_{\alpha\beta} \\ = \frac{\eta}{M^3} j^{\rho\sigma} \left\{ \left(j^{\alpha\beta} (ieF_{\rho\sigma} + \frac{1}{2} j^{\mu\nu} R_{\mu\nu\rho\sigma}) - (ieF_{\rho\sigma} + \frac{1}{2} j^{\mu\nu} R_{\mu\nu\rho\sigma}) j^{\alpha\beta} \right) + M_{\rho\sigma} j^{\alpha\beta} \right\} M_{\alpha\beta} \\ = \frac{\eta}{M^3} j^{\rho\sigma} \left\{ \frac{1}{2} [j^{\alpha\beta}, j^{\mu\nu}]_- R_{\mu\nu\rho\sigma} + M_{\rho\sigma} j^{\alpha\beta} \right\} M_{\alpha\beta}; \quad (130)$$

Thus, we get

$$A = \frac{\eta}{M^3} \left\{ (j^{\rho\sigma} M_{\rho\sigma}) (j^{\alpha\beta} M_{\alpha\beta}) + \frac{1}{2} j^{\rho\sigma} [j^{\alpha\beta}, j^{\mu\nu}]_- M_{\alpha\beta} R_{\mu\nu\rho\sigma} \right\}. \quad (131)$$

Taking in mind the identity [3]

$$j^{\alpha\beta} M_{\alpha\beta} = (ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4}R),$$

we rewrite the last expression as follows

$$A = \frac{\eta}{M^3} j^{\rho\sigma} j^{\alpha\beta} M_{\rho\sigma} M_{\alpha\beta} = \frac{\eta}{M^3} \left\{ \left(ieF_{\rho\sigma} j^{\rho\sigma} - \frac{1}{4}R \right)^2 + \frac{1}{2} j^{\rho\sigma} [j^{\alpha\beta}, j^{\mu\nu}]_- M_{\alpha\beta} R_{\mu\nu\rho\sigma} \right\}. \quad (132)$$

We should detail the second term with commutator

$$\begin{aligned} B &= \frac{1}{2} j^{\rho\sigma} [j^{\alpha\beta}, j^{\mu\nu}]_- R_{\mu\nu\rho\sigma} M_{\alpha\beta} \\ &= \frac{1}{2} j^{\rho\sigma} \left[(g^{\beta\mu} j^{\alpha\nu} - g^{\beta\nu} j^{\alpha\mu}) - (g^{\alpha\mu} j^{\beta\nu} - g^{\alpha\nu} j^{\beta\mu}) \right] R_{\mu\nu\rho\sigma} M_{\alpha\beta} \\ &= \frac{1}{2} j^{\rho\sigma} \left[R_{\nu\rho\sigma}^{\beta} j^{\alpha\nu} - R_{\mu}^{\beta}{}_{\rho\sigma} j^{\alpha\mu} - R_{\nu\rho\sigma}^{\alpha} j^{\beta\nu} + R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\beta\mu} \right] M_{\alpha\beta} \\ &= \frac{1}{2} j^{\rho\sigma} \left[(R_{\nu\rho\sigma}^{\beta} j^{\alpha\nu} M_{\alpha\beta} - R_{\nu\rho\sigma}^{\alpha} j^{\beta\nu} M_{\alpha\beta}) + (R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\beta\mu} M_{\alpha\beta} - R_{\mu}^{\beta}{}_{\rho\sigma} j^{\alpha\mu} M_{\alpha\beta}) \right]. \end{aligned}$$

Making the needed change in mute indices, $\alpha \Rightarrow \beta, \beta \Rightarrow \alpha$, we get

$$\begin{aligned} B &= \frac{1}{2} j^{\rho\sigma} \left[(R_{\nu\rho\sigma}^{\beta} j^{\alpha\nu} M_{\alpha\beta} - R_{\nu\rho\sigma}^{\beta} j^{\alpha\nu} M_{\beta\alpha}) + (R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\beta\mu} M_{\alpha\beta} - R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\beta\mu} M_{\beta\alpha}) \right] \\ &= \frac{1}{2} j^{\rho\sigma} \left[(R_{\nu\rho\sigma}^{\beta} j^{\alpha\nu} M_{\alpha\beta} + R_{\nu\rho\sigma}^{\beta} j^{\alpha\nu} M_{\alpha\beta}) + (R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\beta\mu} M_{\alpha\beta} + R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\beta\mu} M_{\alpha\beta}) \right] \\ &= \frac{1}{2} j^{\rho\sigma} \left[2R_{\nu\rho\sigma}^{\beta} j^{\alpha\nu} M_{\alpha\beta} + 2R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\beta\mu} M_{\alpha\beta} \right] = R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} M_{\alpha\beta} + R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\rho\sigma} j^{\beta\mu} M_{\alpha\beta} = B. \end{aligned}$$

Making the changes, $\mu \Rightarrow \nu, \alpha \Rightarrow \beta, \beta \Rightarrow \alpha$, we get

$$\begin{aligned} B &= R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} M_{\alpha\beta} + R_{\mu}^{\alpha}{}_{\rho\sigma} j^{\rho\sigma} j^{\beta\mu} M_{\alpha\beta} \\ &= R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} M_{\alpha\beta} + R_{\nu}^{\beta}{}_{\rho\sigma} j^{\rho\sigma} j^{\alpha\nu} M_{\beta\alpha} = R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} M_{\alpha\beta} + R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} M_{\alpha\beta} \\ &= 2R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} M_{\alpha\beta} = B. \end{aligned} \quad (133)$$

Allowing for expression for $M_{\alpha\beta}$, we derive

$$\begin{aligned} B &= 2R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} M_{\alpha\beta} = 2R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} \left(ieF_{\alpha\beta} + \frac{1}{2} j^{\epsilon\tau} R_{\epsilon\tau\alpha\beta} \right) \\ &= 2ieF_{\alpha\beta} R_{\nu\rho\sigma}^{\beta} (j^{\rho\sigma} j^{\alpha\nu}) + R_{\nu\rho\sigma}^{\beta} R_{\epsilon\tau\alpha\beta} j^{\rho\sigma} j^{\alpha\nu} j^{\epsilon\tau}. \end{aligned} \quad (134)$$

Let us consider the product of two generators

$$\begin{aligned} j^{\rho\sigma} j^{\alpha\nu} &= \frac{1}{16} (\gamma^{\rho} \gamma^{\sigma} - \gamma^{\sigma} \gamma^{\rho}) (\gamma^{\alpha} \gamma^{\nu} - \gamma^{\nu} \gamma^{\alpha}) \\ &= \frac{1}{16} \left[\gamma^{\rho} (\gamma^{\sigma} \gamma^{\alpha} \gamma^{\nu} - \gamma^{\sigma} \gamma^{\nu} \gamma^{\alpha}) - \gamma^{\sigma} (\gamma^{\rho} \gamma^{\alpha} \gamma^{\nu} - \gamma^{\rho} \gamma^{\nu} \gamma^{\alpha}) \right] \\ &= \frac{1}{16} \left\{ \gamma^{\rho} \left[(\gamma^{\sigma} g^{\alpha\nu} - \gamma^{\alpha} g^{\sigma\nu} + \gamma^{\nu} g^{\sigma\alpha} + i\gamma^5 \epsilon^{\sigma\alpha\nu\delta} \gamma_{\delta}) - (\gamma^{\sigma} g^{\nu\alpha} - \gamma^{\nu} g^{\sigma\alpha} + \gamma^{\alpha} g^{\nu\sigma} + i\gamma^5 \epsilon^{\sigma\nu\alpha\delta} \gamma_{\delta}) \right] \right. \\ &\quad \left. - \gamma^{\sigma} \left[(\gamma^{\rho} g^{\alpha\nu} - \gamma^{\alpha} g^{\rho\nu} + \gamma^{\nu} g^{\rho\alpha} + i\gamma^5 \epsilon^{\rho\alpha\nu\delta} \gamma_{\delta}) - (\gamma^{\rho} g^{\nu\alpha} - \gamma^{\nu} g^{\rho\alpha} + \gamma^{\alpha} g^{\rho\nu} + i\gamma^5 \epsilon^{\rho\nu\alpha\delta} \gamma_{\delta}) \right] \right\} \\ &= \frac{1}{8} \left\{ \gamma^{\rho} \left[-\gamma^{\alpha} g^{\sigma\nu} + \gamma^{\nu} g^{\sigma\alpha} + i\gamma^5 \epsilon^{\sigma\alpha\nu\delta} \gamma_{\delta} \right] - \gamma^{\sigma} \left[-\gamma^{\alpha} g^{\rho\nu} + \gamma^{\nu} g^{\rho\alpha} + i\gamma^5 \epsilon^{\rho\alpha\nu\delta} \gamma_{\delta} \right] \right\} \end{aligned}$$

Therefore, we derive

$$\begin{aligned} C &= ieF_{\alpha\beta} R_{\nu\rho\sigma}^{\beta} j^{\rho\sigma} j^{\alpha\nu} = \frac{ie}{4} F_{\alpha\beta} \left\{ -\gamma^{\rho} \gamma^{\alpha} g^{\sigma\nu} R_{\nu\rho\sigma}^{\beta} + \gamma^{\rho} \gamma^{\nu} g^{\sigma\alpha} R_{\nu\rho\sigma}^{\beta} + i\gamma^{\rho} \gamma^5 \epsilon^{\sigma\alpha\nu\delta} \gamma_{\delta} R_{\nu\rho\sigma}^{\beta} \right. \\ &\quad \left. + \gamma^{\sigma} \gamma^{\alpha} g^{\rho\nu} R_{\nu\rho\sigma}^{\beta} - \gamma^{\sigma} \gamma^{\nu} g^{\rho\alpha} R_{\nu\rho\sigma}^{\beta} - i\gamma^{\sigma} \gamma^5 \epsilon^{\rho\alpha\nu\delta} \gamma_{\delta} R_{\nu\rho\sigma}^{\beta} \right\} \\ &= \frac{ie}{4} F_{\alpha\beta} \left\{ -\gamma^{\rho} \gamma^{\alpha} R_{\nu\rho\sigma}^{\beta} g^{\sigma\nu} + \gamma^{\rho} \gamma^{\nu} R_{\nu\rho\sigma}^{\beta} g^{\sigma\alpha} - i\gamma^{\rho} \gamma_{\delta} \gamma^5 \epsilon^{\sigma\alpha\nu\delta} R_{\nu\rho\sigma}^{\beta} + \gamma^{\sigma} \gamma^{\alpha} R_{\nu\rho\sigma}^{\beta} g^{\rho\nu} - \gamma^{\sigma} \gamma^{\nu} R_{\nu\rho\sigma}^{\beta} g^{\rho\alpha} + i\gamma^{\sigma} \gamma_{\delta} \gamma^5 \epsilon^{\rho\alpha\nu\delta} R_{\nu\rho\sigma}^{\beta} \right\}. \end{aligned}$$

Taking into account the symmetry properties of the curvature tensor, we derive

$$C = \frac{ie}{4} F_{\alpha\beta} \left\{ -\gamma^{\rho} \gamma^{\alpha} R_{\sigma\rho}^{\beta} + (-\gamma^{\nu} \gamma^{\rho} + 2g^{\rho\nu}) R_{\nu\rho}^{\beta} - i\gamma^{\rho} \gamma_{\delta} \gamma^5 \epsilon^{\sigma\alpha\nu\delta} R_{\nu\rho}^{\beta} \right\}$$

$$\begin{aligned}
 & +\gamma^\sigma \gamma^\alpha R^{\beta\rho}_{\rho\sigma} + (\gamma^\nu \gamma^\sigma - 2g^{\nu\sigma}) R^{\beta}_{\nu}{}^\alpha{}_\sigma - i\gamma^\rho \gamma_\delta \gamma^5 \epsilon^{\sigma\alpha\nu\delta} R^{\beta}_{\nu\rho\sigma} \Big\} \\
 & = \frac{ie}{4} F_{\alpha\beta} \Big\{ -\gamma^\rho \gamma^\alpha R^{\beta}_{\rho}{}^\alpha + (-\gamma^\nu \gamma^\rho R^{\beta}_{\nu\rho}{}^\alpha - 2R^{\alpha\beta}) - 2i\gamma^\rho \gamma_\delta \gamma^5 \epsilon^{\sigma\alpha\nu\delta} R^{\beta}_{\nu\rho\sigma} - \gamma^\rho \gamma^\alpha R^{\beta}_{\rho}{}^\alpha + (-\gamma^\nu \gamma^\rho R^{\beta}_{\nu\rho}{}^\alpha + 2R^{\alpha\beta}) \Big\}, \\
 & \text{that is}
 \end{aligned}$$

$$C = \frac{ie}{2} F_{\alpha\beta} \Big\{ -\gamma^\rho \gamma^\alpha R^{\beta}_{\rho}{}^\alpha - \gamma^\nu \gamma^\rho R^{\beta}_{\nu\rho}{}^\alpha - i\gamma^\rho \gamma_\delta \gamma^5 \epsilon^{\sigma\alpha\nu\delta} R^{\beta}_{\nu\rho\sigma} \Big\}.$$

Using the identities

$$\gamma^\rho \gamma^\alpha = (2j^{\rho\alpha} + g^{\rho\alpha}), \quad \gamma^\nu \gamma^\rho = (2j^{\nu\rho} + g^{\nu\rho}), \quad \gamma^\rho \gamma_\delta = (2j^\rho_\delta + \delta^\rho_\delta),$$

we arrive at

$$\begin{aligned}
 C &= \frac{ie}{2} F_{\alpha\beta} \Big\{ -(2j^{\rho\alpha} R^{\beta}_{\rho}{}^\alpha + g^{\rho\alpha} R^{\beta}_{\rho}{}^\alpha) - (2j^{\nu\rho} R^{\beta}_{\nu\rho}{}^\alpha + g^{\nu\rho} R^{\beta}_{\nu\rho}{}^\alpha) - i(2j^\rho_\delta + \delta^\rho_\delta) \gamma^5 \epsilon^{\sigma\alpha\nu\delta} R^{\beta}_{\nu\rho\sigma} \Big\} \\
 &= \frac{ie}{2} F_{\alpha\beta} \Big\{ -2j^{\rho\alpha} R^{\beta}_{\rho}{}^\alpha - R^{\alpha\beta} - 2j^{\nu\rho} R^{\beta}_{\nu\rho}{}^\alpha + R^{\alpha\beta} - 2ij^\rho_\delta \gamma^5 \epsilon^{\sigma\alpha\nu\delta} R^{\beta}_{\nu\rho\sigma} \Big\} \\
 &= \frac{ie}{2} F_{\alpha\beta} \Big\{ -2j^{\rho\alpha} R^{\beta}_{\rho}{}^\alpha - 2j^{\nu\rho} R^{\beta}_{\nu\rho}{}^\alpha - 2i\gamma^5 j^{\rho\delta} \epsilon^{\sigma\alpha\nu\delta} R^{\beta}_{\nu\rho\sigma} \Big\} = C.
 \end{aligned} \tag{135}$$

Collecting results together, we get

$$\begin{aligned}
 & \frac{\eta}{M^3} j^{\rho\sigma} j^{\alpha\beta} M_{\rho\sigma} M_{\alpha\beta} = \frac{\eta}{M^3} \left[\left(ieF_{\rho\sigma} j^{\rho\sigma} - \frac{1}{4} R \right)^2 \right. \\
 & \left. + \frac{ie}{2} F_{\alpha\beta} \left(-2j^{\rho\alpha} R^{\beta}_{\rho}{}^\alpha - 2j^{\nu\rho} R^{\beta}_{\nu\rho}{}^\alpha - 2i\gamma^5 j^{\rho\delta} \epsilon^{\sigma\alpha\nu\delta} R^{\beta}_{\nu\rho\sigma} \right) + R^{\beta}_{\nu\rho\sigma} R_{\epsilon\tau\alpha\beta} j^{\rho\sigma} j^{\alpha\nu} j^{\epsilon\tau} \right].
 \end{aligned} \tag{136}$$

Therefore, the basic equation takes on the form

$$\begin{aligned}
 & \left\{ (\gamma^\sigma D_\sigma + M) + \frac{\mu}{M} \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R \right) + \frac{\sigma}{M^2} (\gamma^\rho D_\rho) \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R \right) \right. \\
 & \left. + \frac{\eta}{M^3} \left[D^\sigma D_\sigma \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R \right) + \left(ieF_{\rho\sigma} j^{\rho\sigma} - \frac{1}{4} R \right)^2 \right. \right. \\
 & \left. \left. + \frac{ie}{2} F_{\alpha\beta} \left(2j^{\alpha\rho} R_{\rho}{}^\beta - 2j^{\nu\rho} R^{\beta}_{\nu\rho}{}^\alpha - 2i\gamma^5 j^{\rho\delta} \epsilon_{\delta}{}^{\alpha\nu\sigma} R^{\beta}_{\nu\rho\sigma} \right) + R^{\beta}_{\nu\rho\sigma} R_{\epsilon\tau\alpha\beta} j^{\rho\sigma} j^{\alpha\nu} j^{\epsilon\tau} \right] \right\} \Psi = 0.
 \end{aligned} \tag{137}$$

The operator $D^\sigma D_\sigma$ is detailed as follows

$$\begin{aligned}
 D^\sigma D_\sigma &= (\nabla^\sigma + \Gamma^\sigma)(\nabla_\sigma + \Gamma_\sigma) = \nabla^\sigma \nabla_\sigma + (\nabla_\sigma \Gamma^\sigma) + \Gamma^\sigma \nabla_\sigma + \Gamma^\sigma \nabla_\sigma + \Gamma^\sigma \Gamma_\sigma \\
 &= \nabla^\sigma \nabla_\sigma + (\gamma_\sigma \Gamma^\sigma - \Gamma^\sigma \gamma_\sigma) + 2\Gamma^\sigma \nabla_\sigma + \Gamma^\sigma \Gamma_\sigma.
 \end{aligned}$$

The last term in (137) may be presented differently

$$R^{\beta}_{\nu\rho\sigma} R_{\epsilon\tau\alpha\beta} j^{\rho\sigma} j^{\alpha\nu} j^{\epsilon\tau} \equiv -(R^{\beta}_{\nu\rho\sigma} j^{\rho\sigma}) j^{\nu\alpha} (R_{\alpha\beta\epsilon\tau} j^{\epsilon\tau}) \equiv -X^{\beta}_{\nu} j^{\nu\alpha} X_{\alpha\beta}. \tag{138}$$

It is convenient to transform the last term in (137) to tetrad form:

$$R^{\beta}_{\nu\rho\sigma} R_{\epsilon\tau\alpha\beta} j^{\rho\sigma} j^{\alpha\nu} j^{\epsilon\tau} = -\left(R^d{}_{nab}(x) j^{ab} \right) j^{nm} \left(R_{mdkl}(x) j^{kl} \right).$$

Therefore, the basic equation takes on the form

$$\begin{aligned}
 & \left\{ \left(\gamma^\sigma D_\sigma + M \right) + \frac{\mu}{M} \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R \right) + \frac{\sigma}{M^2} (\gamma^\rho D_\rho) \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R \right) \right. \\
 & \left. + \frac{\eta}{M^3} \left[D^\sigma D_\sigma \left(ieF_{\alpha\beta} j^{\alpha\beta} - \frac{1}{4} R \right) + \left(ieF_{\rho\sigma} j^{\rho\sigma} - \frac{1}{4} R \right)^2 + \right. \right. \\
 & \left. \left. + \frac{ie}{2} F_{\alpha\beta} \left(2j^{\alpha\rho} R_{\rho}{}^\beta - 2j^{\nu\rho} R^{\beta}_{\nu\rho}{}^\alpha - 2i\gamma^5 j^{\rho\delta} \epsilon_{\delta}{}^{\alpha\nu\sigma} R^{\beta}_{\nu\rho\sigma} \right) \right. \right. \\
 & \left. \left. - \left(R^d{}_{nab}(x) j^{ab} \right) j^{nm} \left(R_{mdkl}(x) j^{kl} \right) \right] \right\} \Psi = 0.
 \end{aligned} \tag{139}$$

Evidently, it is possible to detail additionally the term with the product of two Riemann tensors.

Conclusions

As seen, in the derived extended equation for spin $1/2$ particle, there arises a number of additional geometrical interaction terms through the Ricci scalar $R(x)$ and tensor R_{ab} , and the Riemann curvature tensor $R_{mnkl}(x)$. The contribution of Ricci tensor and Riemann curvature tensor differ from zero only if the third parameter η vanishes.

References

- [1] Pletyukhov, V.A. Relativistic wave equations and internal degrees of freedom / V.A. Pletyukhov, V. M. Red'kov, V. I. Strazhev // Belarussian Science, Minsk. – 2015. – 328 p.
- [2] Elementary Particles with Internal Structure in External Fields. Vol I. General Theory; Vol II. Physical Problems / V.V. Kisel, E.M. Ovsyuk, O.V. Veko, Y.A. Voynova, V. Balan, V.M. Red'kov. – New York: Nova Science Publishers Inc., 2018. – 402 p., 404 p.
- [3] Red'kov V.M. Field particles in Riemannian space and the Lorentz group. – Belarussian Science, Minsk, 2009. 486 p. (in Russian).